# Complexity in Multi-Delay Physiological Feedback Systems

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# **OUTLINE**

## 1) Chaos to simple dynamics via many delays: how?

→ COMPLEXITY IS NOT NECESSARILY HIGH WITH MULTIPLE DELAYS

→ COMPLEXITY COLLAPSES CAN HAPPEN WITH EVEN A FEW DELAYS

2) Complexity collapse in neural nets?

→TRANSIENT CHAOS, RANDOM PERIODS OF SYNCHRONY

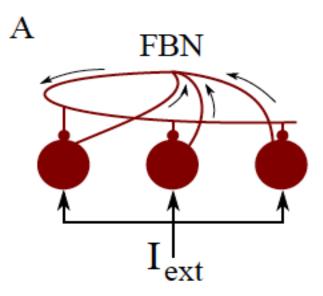
SIMPLE PUNCHLINES

• More delays does not imply more complexity

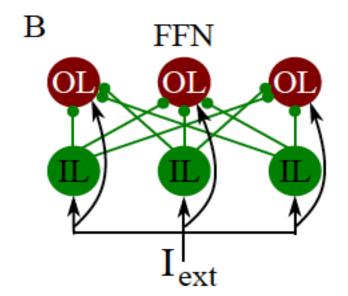
 Delayed feedback loops in physiology may be able to do novel forms of:

- Random "number" generation
- Prediction
- Deep learning

FB: DELAYED FEEDBACKFF: FEEDFORWARD+DELAY

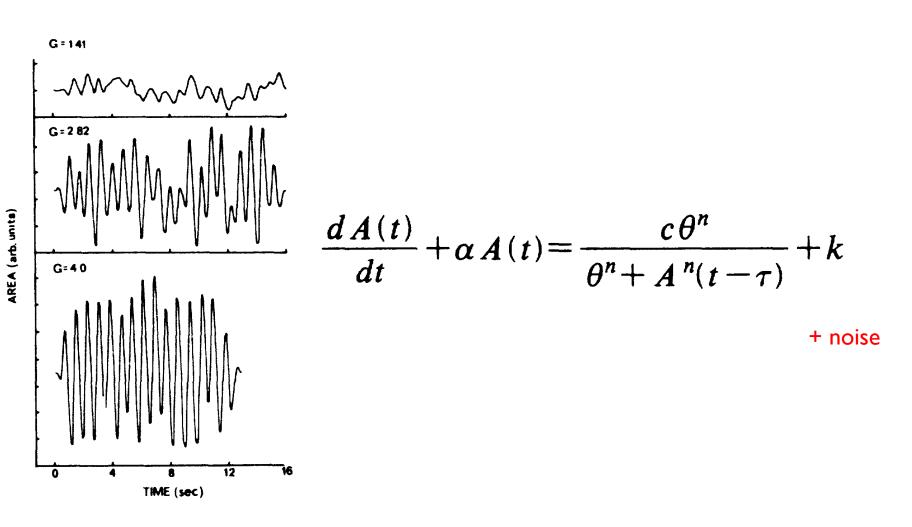


Feedback delay (INHIB)



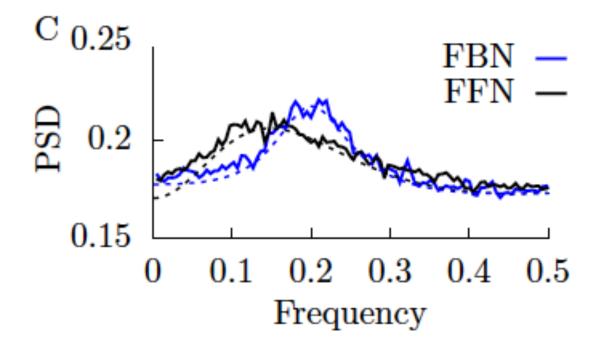
Feedforward delay

#### (HUMAN) PUPIL AREA OSCILLATIONS BY CONTROLLING THE FEEDBACK GAIN



Longtin, Milton, Bos, Mackey, Phys. Rev. A. 1990

### BOTH DISPLAY A PEAK !



Many ways to get rhythms!!

Payeur & Longtin, Phys. Rev. E. 2015

- Increasing a single delay increases entropy (Farmer, Physica D 1984)
- But what about the number of delays?

#### Increased complexity

- Fisher et al. PRL 1994: add 2nd delay to laser system  $\rightarrow$  *chaos to hyperchaos*
- Xu et al., Optics Lett. 2017: large number of random delays → *hyperchaos for random number generation*

#### Decreased complexity

- Ahlborn+Parlitz, PRL2004: 2<sup>nd</sup> delay on Chua circuit: chaos replaced by limit cycles
- Mensour+Longtin, PhysLett A 1995: controlling chaos in Mackey-Glass with a 2<sup>nd</sup> delay

# PARADOX

1) Multiple delays can increase the complexity (entropy)

2) But distributed delays have relatively lower complexity

What happens in between??

## SEMICONDUCTOR LASER MODEL WITH OPTICAL FEEDBACK

• Semiconductor laser with multiple optical feedbacks can be modeled with Lang-Kobayashi (LK) equations:

$$\frac{dE(t)}{dt} = (1+i\alpha) \left[ \frac{G_E(N(t)-N_0)}{1+\epsilon|E(t)|^2} - \gamma_E \right] E(t) + \frac{\kappa}{M} \sum_{i=1}^M E(t-\tau_i) e^{-i\omega\tau_i}$$
$$\frac{dN(t)}{dt} = \gamma_N (J_r N_{th} - N(t)) - \frac{G_N(N(t)-N_0)}{1+\epsilon|E(t)|^2} |E(t)|^2$$

> Dynamics of LK model is governed by two main time scales:

#### Average delay and response time

• In distributed delay case, the state of the system depends on a continuum of past states:

$$\frac{dE}{dt} = \int_0^\infty g(\tau) F[N(t), E(t), E(t - \tau)] d\tau,$$
$$\frac{dN(t)}{dt} = \gamma_N [J_r N_{th} - N(t)] - \frac{G_N [N(t) - N_0]}{1 + \epsilon |E(t)|^2} |E(t)|^2$$

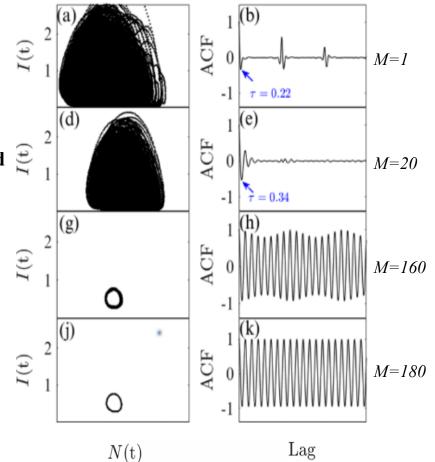
## SEMICONDUCTOR LASER MODEL WITH OPTICAL FEEDBACK

- Weak or moderate degree of chaos: Dominant time scale is the **time delay**
- High degree of chaos: Dominant time scale is the **relaxation oscillation period**
- Torus behavior:

Both relaxation oscillation period and average of the time delays govern the dynamic

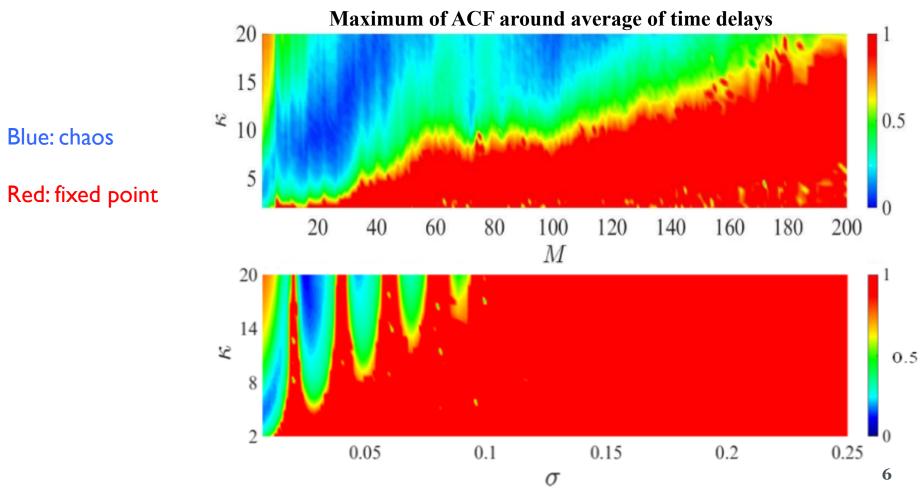
• Limit cycle:

Simple oscillation with relaxation oscillation period



## MIXED EFFECTS WHEN ADDING DELAYS

• Degree of complexity via the maximum of the *autocorrelation function* (ACF) around the time delay.



S. K. Tavakoli and A. Longtin, "Multi-delay complexity collapse," Phys. Rev. Research 2, 033485 (2020)

## FIXED AVERAGE DELAY: M INDUCES BIFURCATIONS

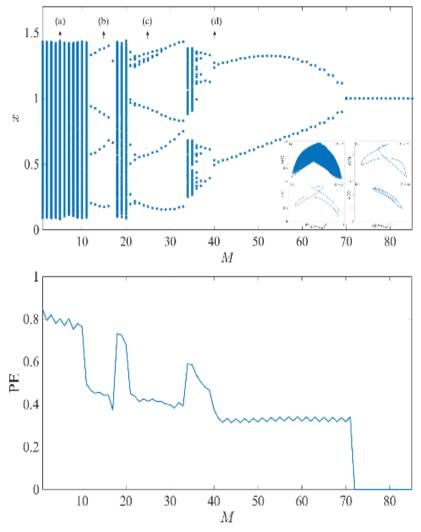
• Mackey Glass equation (standard parameters):

$$\frac{dx}{dt} = -x(t) + \frac{b}{M} \sum_{i=1}^{M} \frac{x(t-\tau_i)}{1+x(t-\tau_i^{10})}$$

• Delays are picked as:

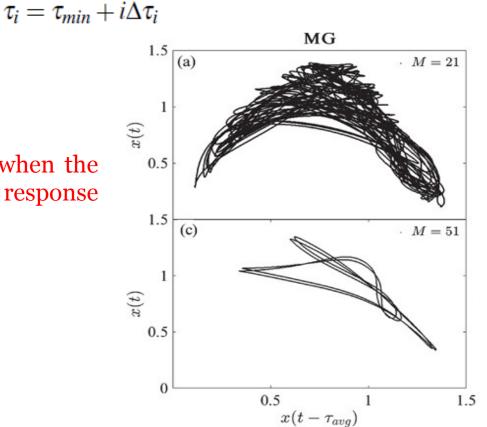
 $au_i = au_{avg} - (i-1)/2\Delta au$ , *i* odd  $au_i = au_{avg} + (i)/2\Delta au$ , *i* even

• Increasingly large periodic windows occur for increasing number of delays *M*, culminating in an inverse period-doubling cascade that ends with a fixed point.

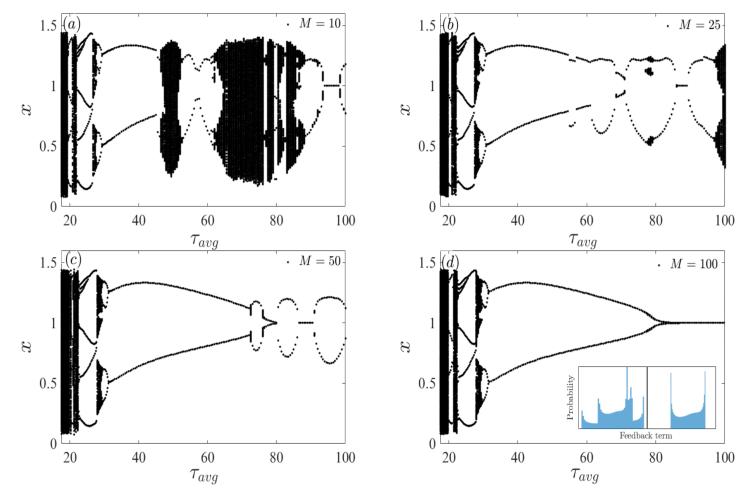


- Whenever the delay-to-response time ratio is large, first order nonlinear DDE's exhibit high-dimensional chaos and multistability.
- We use another scheme to pick delays by keeping the minimum delay fixed and increasing the average of the delays:

Complexity collapse still observed when the time delays are larger than the response time.

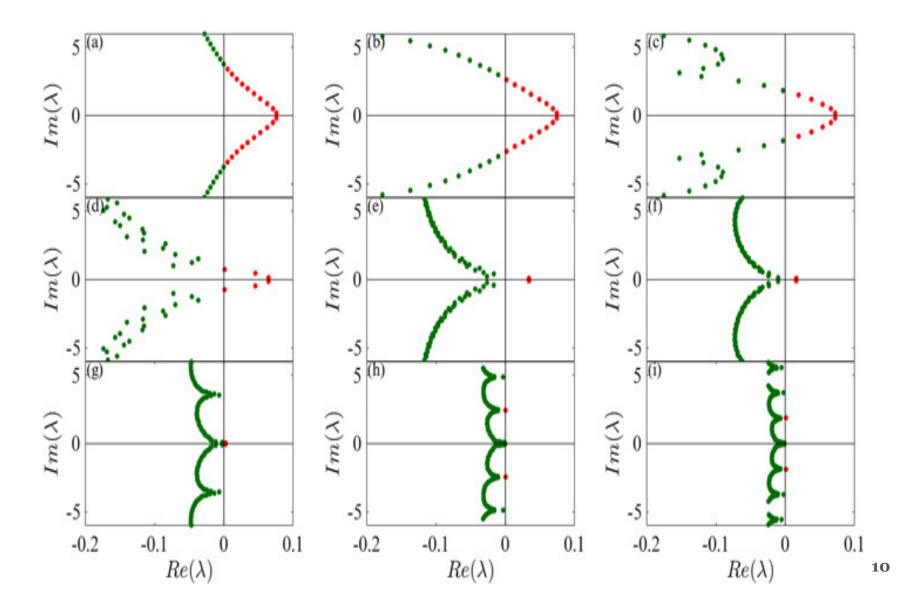


<u>S. Kamyar Tavakoli</u> and <u>André Longtin</u>, "Dynamical invariants and inverse period-doubling cascades 7 in multi-delay systems", Chaos 31, 103129 (2021) • We found that complexity collapse occurs for even few number of delays, provided there is a sufficiently large minimum delay.



• A larger number of delays favors fixed point behavior rather than limit cycles.

Behavior of eigenvalues around fixed point when M=50, for increasing average delay (a) → (i)



• Characteristic equations:

$$\Delta(\lambda) = -\lambda - 1 + \frac{F'(x^*)}{M} \sum_{j=1}^{M} e^{-\lambda \tau_j}$$
$$\lambda = \mu + i\omega$$

• Separating real and imaginary part of the characteristic equation:

$$\frac{\mu + 1}{F'(x^*)} = \frac{e^{-\mu\tau_{min}}}{M} \sum_{i=1}^{M} e^{-\mu(j-1)\Delta\tau} \cos(\omega(\tau_{min} + (j-1)\Delta\tau))$$
$$-\frac{\omega}{F'(x^*)} = \frac{e^{-\mu\tau_{min}}}{M} \sum_{j=1}^{M} e^{-\mu(j-1)\Delta\tau} \sin(\omega(\tau_{min} + (j-1)\Delta\tau))$$

• Putting the above equations in integral form:

$$\frac{\mu+1}{F'(x^*)} \approx \frac{e^{-\mu\tau_{min}}}{T_M + \Delta\tau} \int_0^{T_M} e^{-\mu\tau'} \cos(\omega(\tau_{min} + \tau')) d\tau'$$
$$-\frac{\omega}{F'(x^*)} \approx \frac{e^{-\mu\tau_{min}}}{T_M + \Delta\tau} \int_0^{T_M} e^{-\mu\tau'} \sin(\omega(\tau_{min} + \tau')) d\tau'$$
$$T_m = \tau_{max} \cdot \tau_{min}$$
$$\frac{\mu+1}{F'(x^*)} - \frac{e^{-\mu\tau_{min}} \left[ e^{-\mu T_M} \sin(\omega(\tau_{min} + T_M)) - \sin(\omega\tau_{min}) \right]}{\omega T_M} = 0$$
$$\frac{-\omega}{F'(x^*)} - \frac{e^{-\mu\tau_{min}} \left[ \cos(\omega\tau_{min}) - e^{-\mu T_m} \cos(\omega(\tau_{min} + T_m)) \right]}{\omega T_M} = 0$$

## LYAPUNOV EXPONENTS

- Lyapunov exponents are calculated using Chebyshev polynomial nodes.
- There are  $m_i$ +1 Chebyshev nodes between two-time delays:

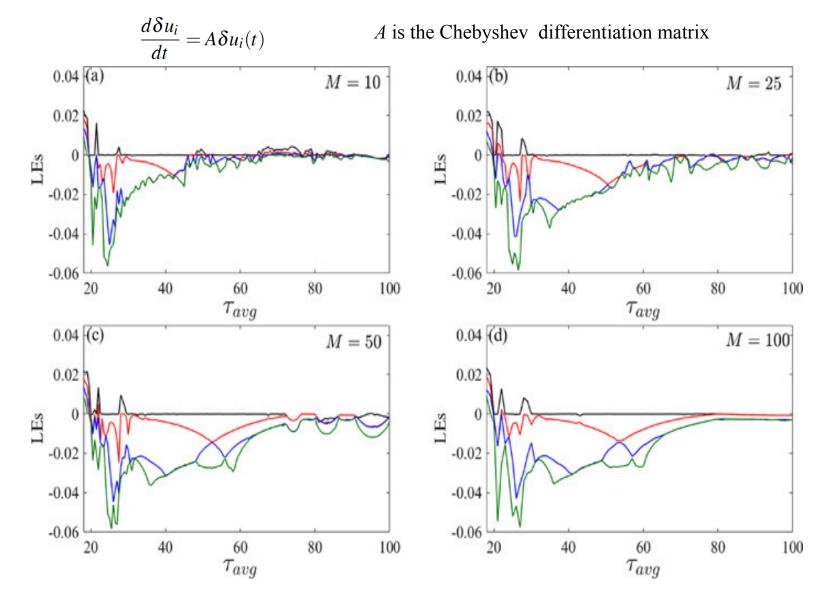
$$\theta_{ij} = \frac{(\tau_i - \tau_{i-1})}{2} \cos\left(\frac{j\pi}{m_i}\right) - \frac{(\tau_i + \tau_{i-1})}{2},$$
  
 $i = 1, \dots, M, \quad j = 0, \dots, m_i,$ 

• Dynamical Equations to estimate Lyapunov spectrum:

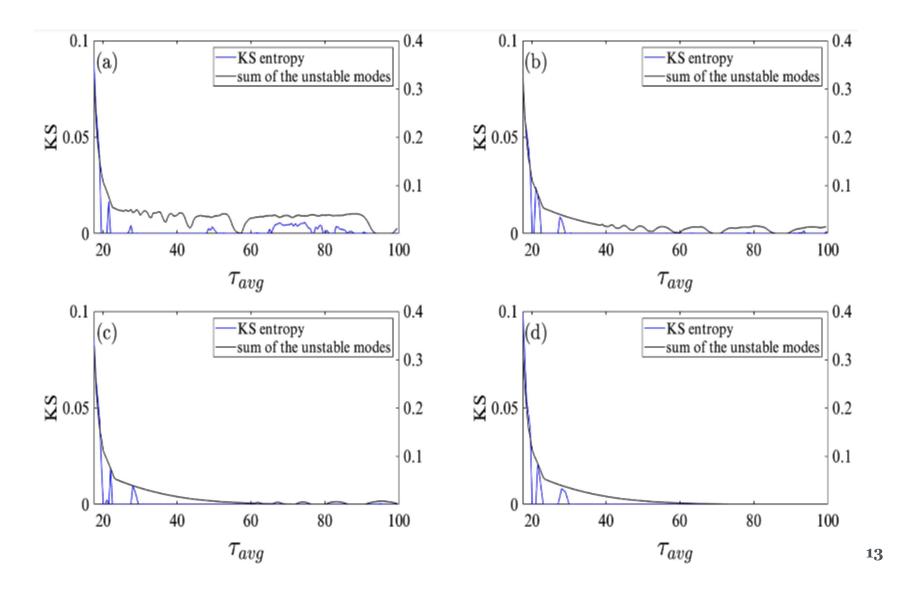
$$A = \begin{pmatrix} \frac{df}{dx} & \cdots & \frac{df}{dx_{\tau_1}} & 0 & \cdots & 0 & \cdots & \frac{df}{dx_{\tau_M}} \\ \frac{df}{dx} & \cdots & \frac{df}{dx_{\tau_1}} & 0 & \cdots & 0 & \cdots & \frac{df}{dx_{\tau_M}} \\ \frac{df}{dx_0} & \cdots & d_{1,m_1} & 0 & \cdots & \vdots & \cdots & 0 \\ \vdots & \ddots & \vdots & \vdots & \cdots & 0 & \cdots & \vdots \\ \frac{dm_{1,0}}{dm_{1,0}} & \cdots & \frac{dm_{1+1,m_1}}{dm_{1+1,m_1}} & 0 & \cdots & 0 \\ \vdots & \ddots & \vdots & \vdots & \cdots & 0 & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots & \cdots & 0 \\ 0 & \cdots & \frac{dm_{1+1,m_1}}{dm_{1+m_2,m_1}} & \cdots & \frac{dm_{1+1,m_1+m_2}}{dm_{1+m_2,m_1+m_2}} & 0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots & \cdots & 0 \\ 0 & \cdots & \frac{dm_{1+m_2,m_1}}{dm_{1+m_2,m_1+m_2}} & 0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots & \cdots & 0 \\ 0 & 0 & \frac{dm_{tor} - m_M + 1, m_{tor} - m_M}{dm_{tor} - m_M - m_M - m_M + 1, m_{tor}} \end{pmatrix}$$

#### Generalization of Breda and Schiava, DCDS2018

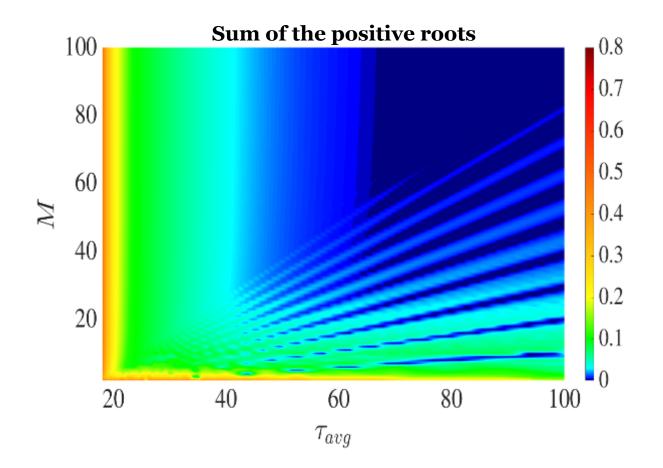
• Lyapunov exponents calculated by integrating the linearized dynamics for the  $u_k$ 



• Comparison between Sum of Positive Lyapunov Exponents and Sum of positive part of eigenvalues



• Sum of Real(eigenvalues): proxy for solution complexity

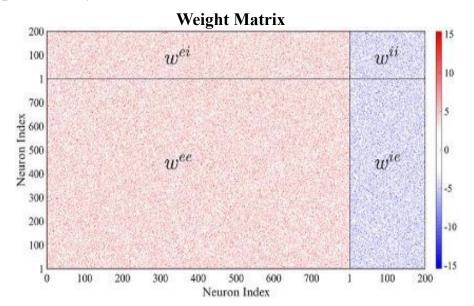


#### \* MODEL MIMICKING SPARSELY CONNECTED BRAIN CIRCUITS

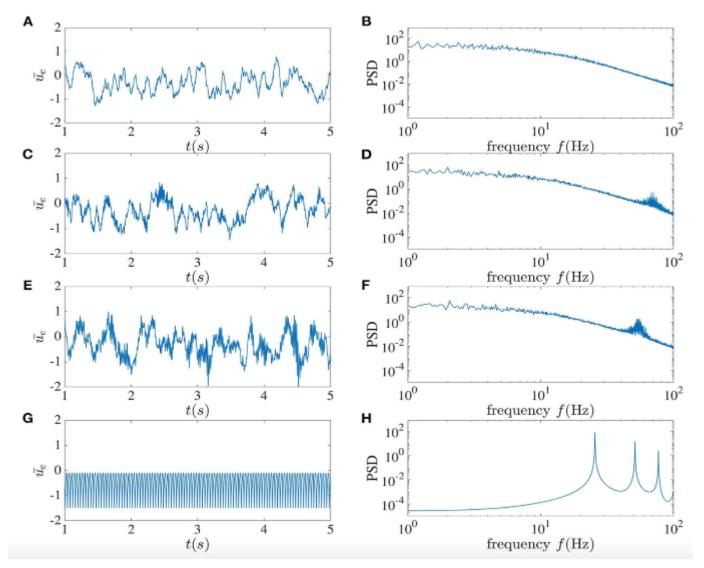
• Recurrent Neural network with multiple local time delays:

$$\begin{aligned} \alpha_e^{-1} \frac{du_j}{dt} &= -u_j + \frac{1}{n_e} \frac{1}{M} \sum_{l=1}^M \sum_{k=1}^{N_e} w_{jk}^{ee} \phi(u_k(t-\tau_l)) + \frac{1}{n_i} \frac{1}{M} \sum_{l=1}^M \sum_{k=1}^{N_i} w_{jk}^{ie} \phi(v_k(t-\tau_l)) \\ \alpha_i^{-1} \frac{dv_j}{dt} &= -v_j + \frac{1}{n_e} \frac{1}{M} \sum_{l=1}^M \sum_{k=1}^{N_e} w_{jk}^{ei} \phi(u_k(t-\tau_l)) + \frac{1}{n_i} \frac{1}{M} \sum_{l=1}^M \sum_{k=1}^{N_i} w_{jk}^{ii} \phi(v_k(t-\tau_l)) \end{aligned}$$

- Rich S, Hutt A, Skinner FK, Valiante TA, Lefebvre J.. Sci Rep. (2020); 10 (1):1–17
- Park, S. H., Griffiths, J. D., Longtin, A., and Lefebvre, J. (2018). Front. Appl. Math.
- We consider a recurrent local network of 80% excitatory and 20% inhibitory rate model neurons with 10% connection probability.

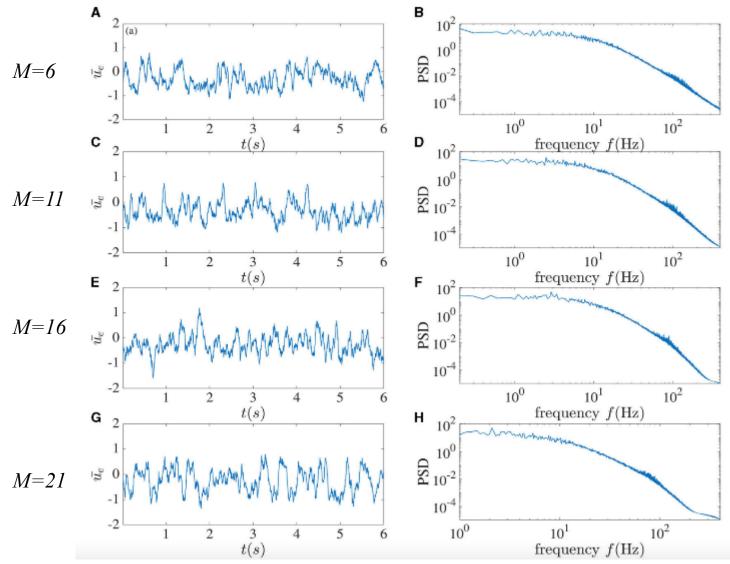


• Recurrent Neural network with single local time delay: Effect of increasing the delay from 2 ms to 10 ms

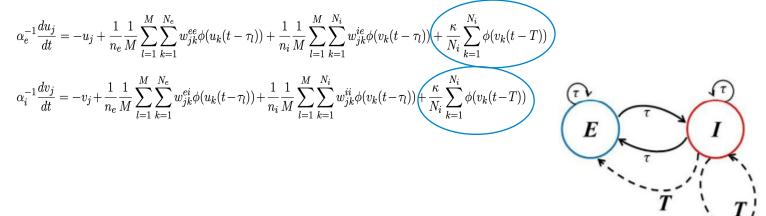


<u>S. Kamyar Tavakoli</u> and <u>André Longtin</u>, "Complexity Collapse, Fluctuating Synchrony, and Transient Chaos in Neural Networks With Delay Clusters", Frontiers Syst. Neurosci. 31, 103129 (2021)

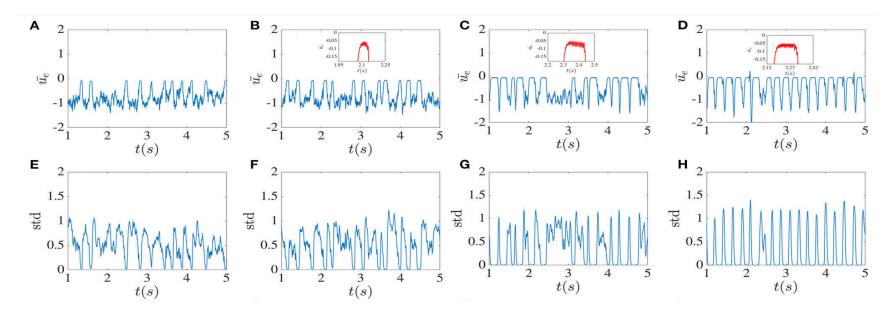
• Recurrent Neural network with local multiple time delays: no complexity collapse!



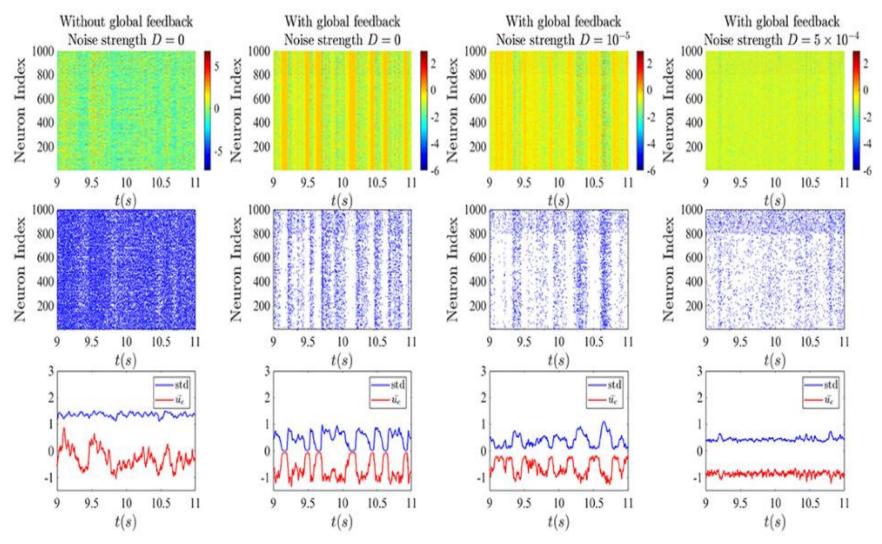
• Recurrent Neural network with global delayed inhibitory global feedback: collapse!



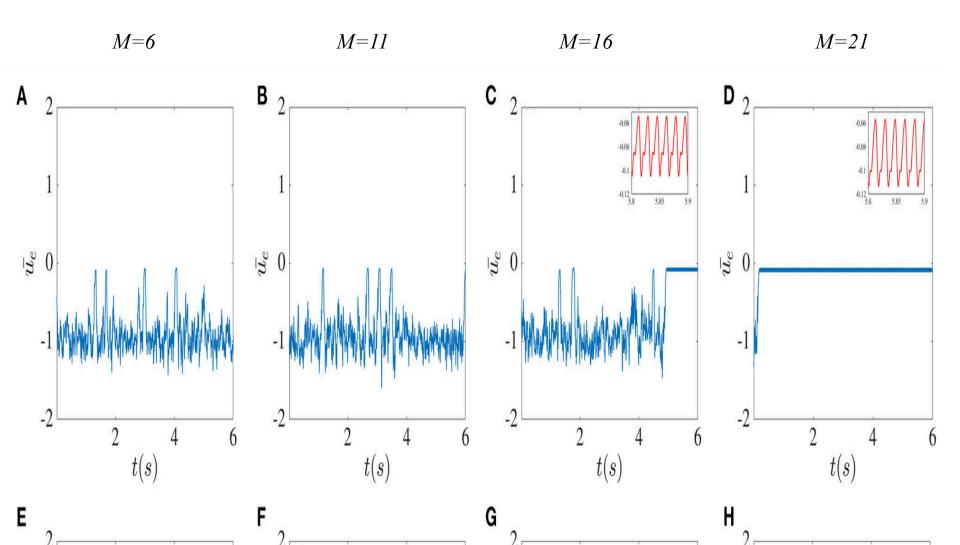
Effect of different global time delays T=5,10,20 and 30 ms



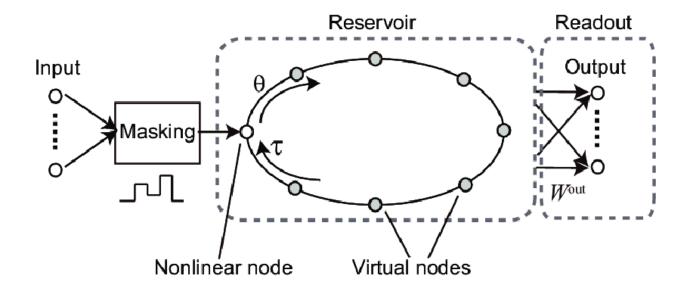
### NEURON ACTIVITIES: STRANGE SYNCHED TEMPORAL RANDOMNESS



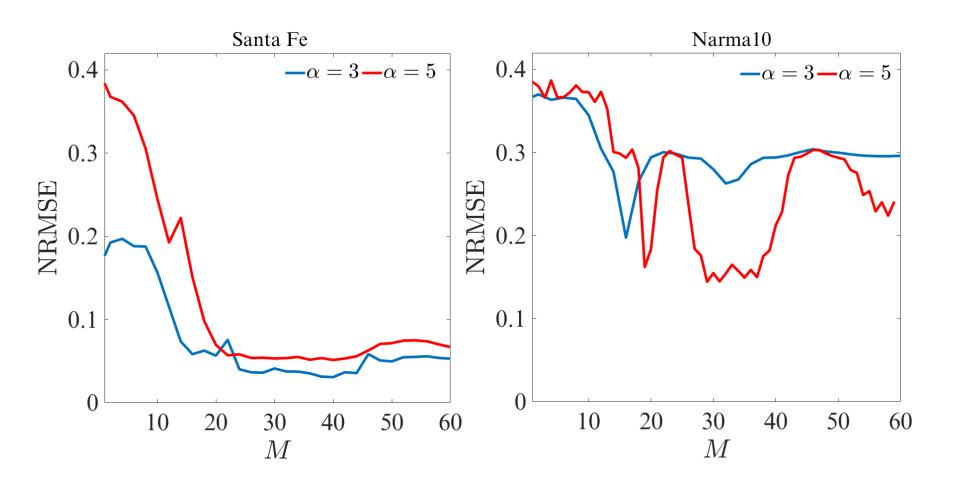
Multiple local delays in the presence of a moderate global delay: transient chaos  $\rightarrow$  collapse



#### RESERVOIR COMPUTING WITH DELAYED LOOPS (APPELTANT ET AL. 2011)







# **\*SUMMARY**

- Transition to simplicity can be abrupt or follow inverse period-doubling sequence in Mackey-Glass blood cell control model
- Large number of delays generally favors simplified dynamics in delayed differential equations in 1-3 variables.
- Chaotic recurrent neural networks may not be simplified by local delays alone.
- The presence of both local delays and global inhibitory feedback may cause collapse to simple dynamics.
- Multiple (but not too many) delays increases entropy: randomness generation
- Multiple delays can benefit reservoir computing:

 $\rightarrow$  proximity of equilibrium and stronger hyper-chaos

#### THANK YOU!