

Explosive synchronization in complex networks

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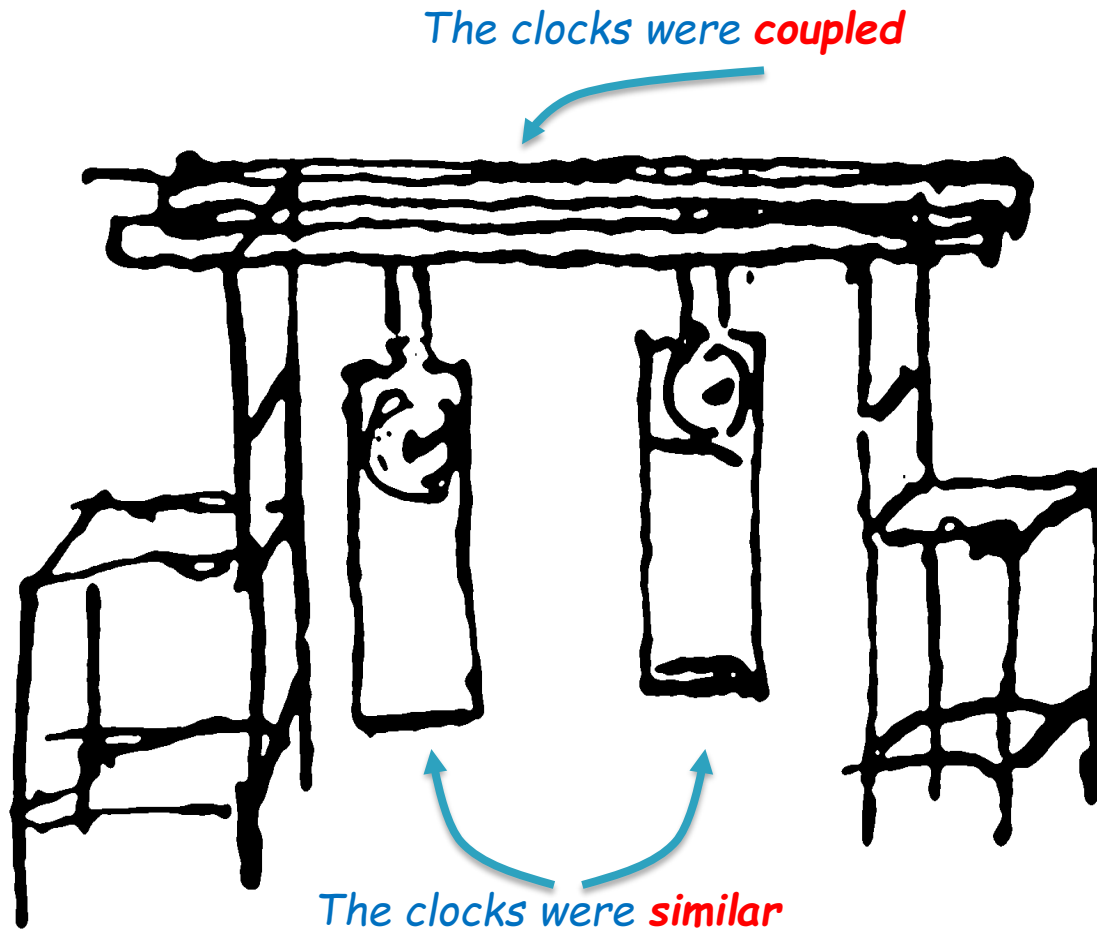
Explosive

Synchronization

in

Complex Networks

The sympathetic clocks of Huyghens

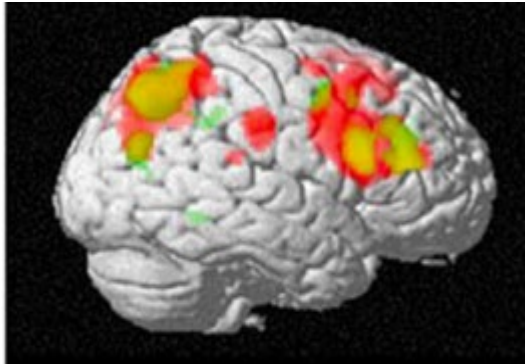


Christiaan Huyghens (1629-1695) discovered what he called "an odd kind of sympathy" between the clocks: regardless of their initial state, both adopted the same rhythm.

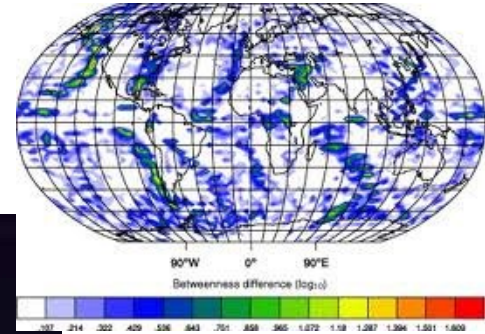
Huygens correctly attributed the synchrony to *tiny forces transmitted by the wooden beam* from which they were suspended.

OK, things synchronize. So what?

Synchrony happens to be the main mechanism for regulating the dynamics and transmit information **in natural ensembles...**



Brain dynamics



World clima ?

Heart beating



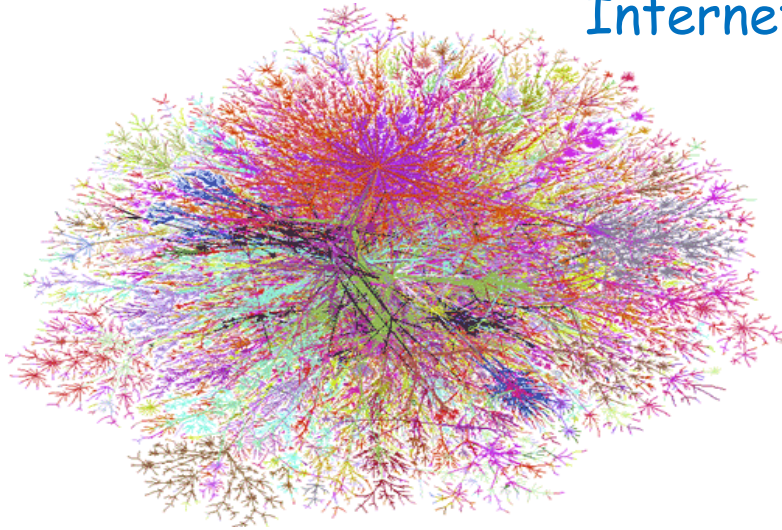
Animal behaviour



Things that synchronize

...and in social or artificial ones

Internet

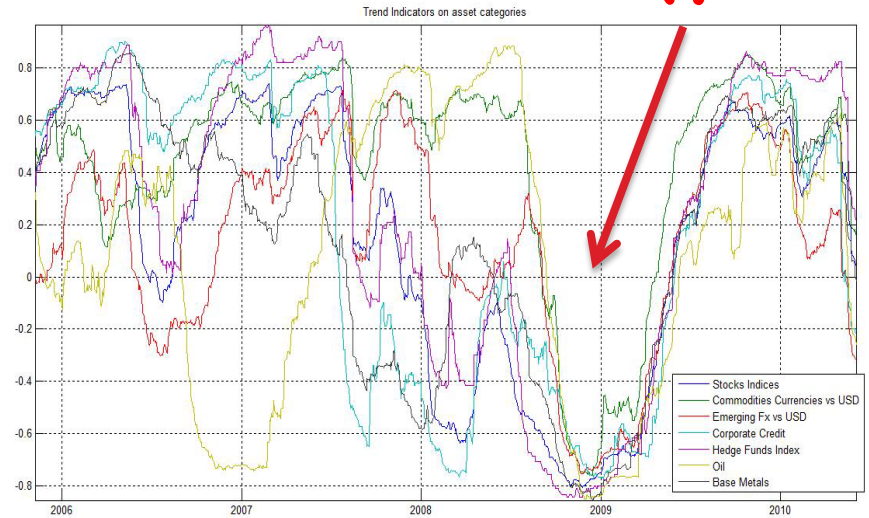


Human behaviour



Financial markets

!?



Power grids



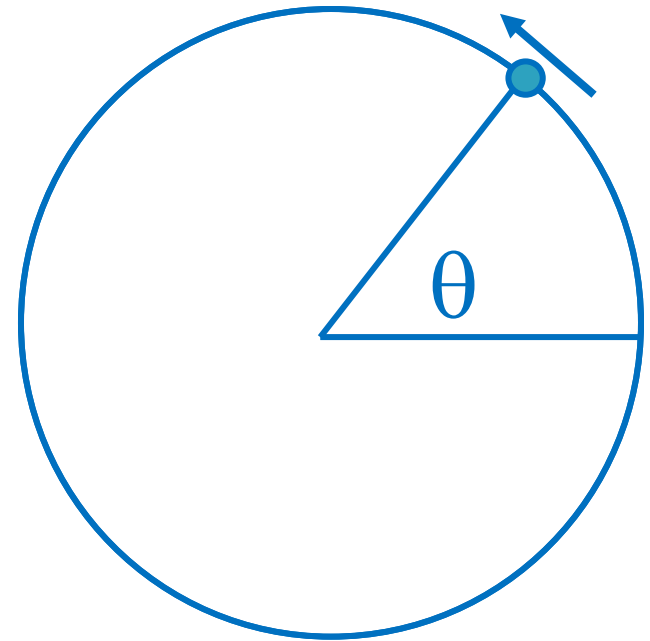
- A periodic oscillator with an intrinsic (or natural) **frequency** ω_n .

- The evolution of each oscillator n is described only by its **phase** θ_n such that

$$\dot{\theta}_n = \omega_n$$

- We are interested in **heterogeneous ensembles**, so we assume the ω_n **frequencies are different**, randomly picked from an (usually known) distribution $g(\omega)$

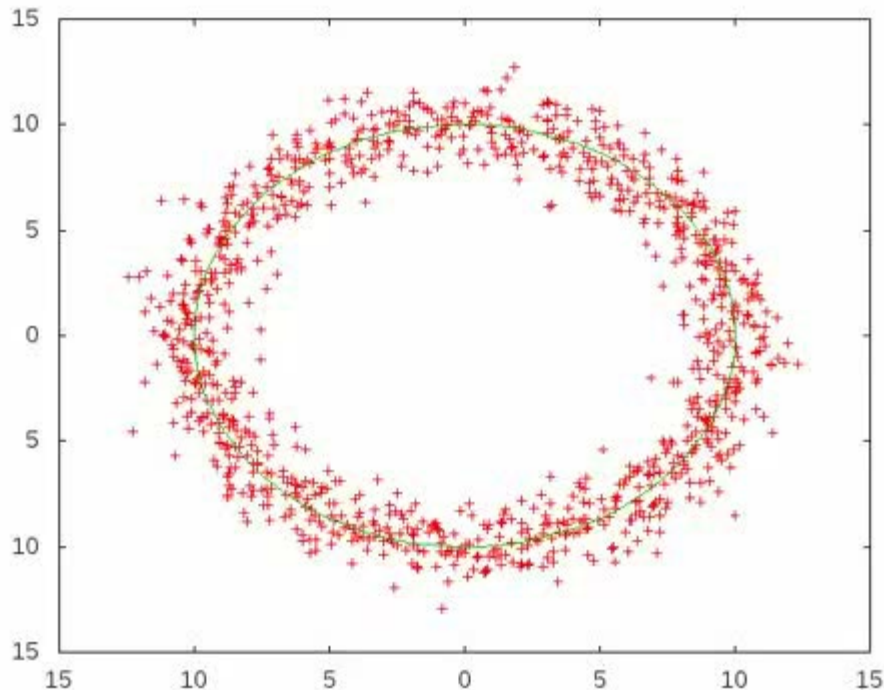
- A large ensemble of N oscillators



Kuramoto ensemble: all-to-all coupling

Sinusoidal *all-to-all* coupling.

$$\dot{\theta}_n = \omega_n$$

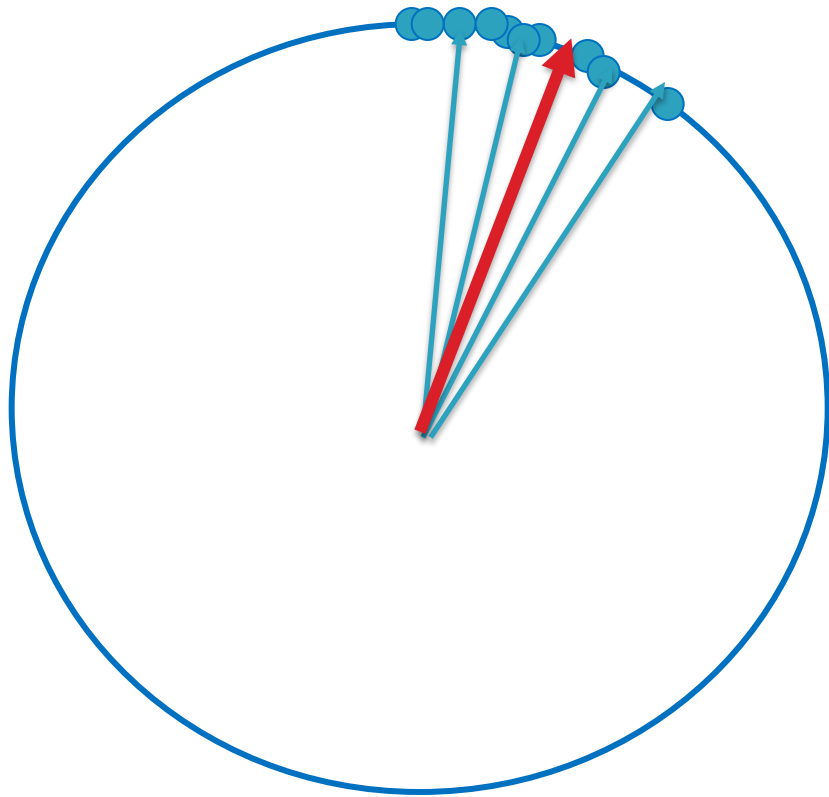


If d is high enough,
phases lock and oscillators
frequency converge
to the average $\langle \omega_n \rangle$

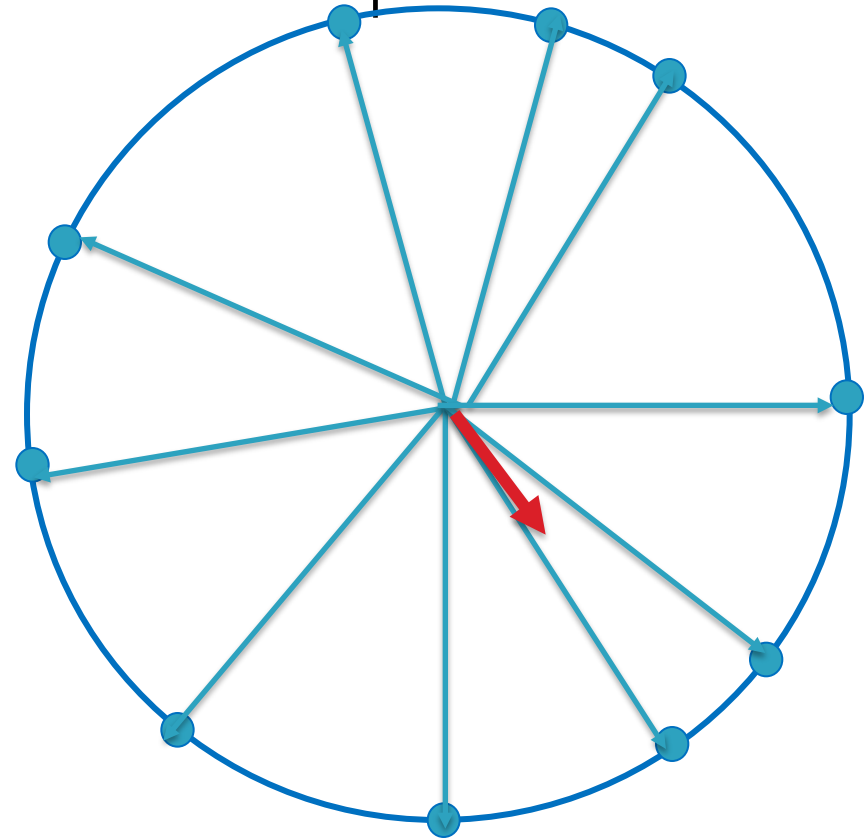
($N=1000$)

Measuring synchronization: Kuramoto order parameter

$$r = \frac{1}{N} \left| \sum_{m=1}^N \exp(i\theta_m) \right|$$

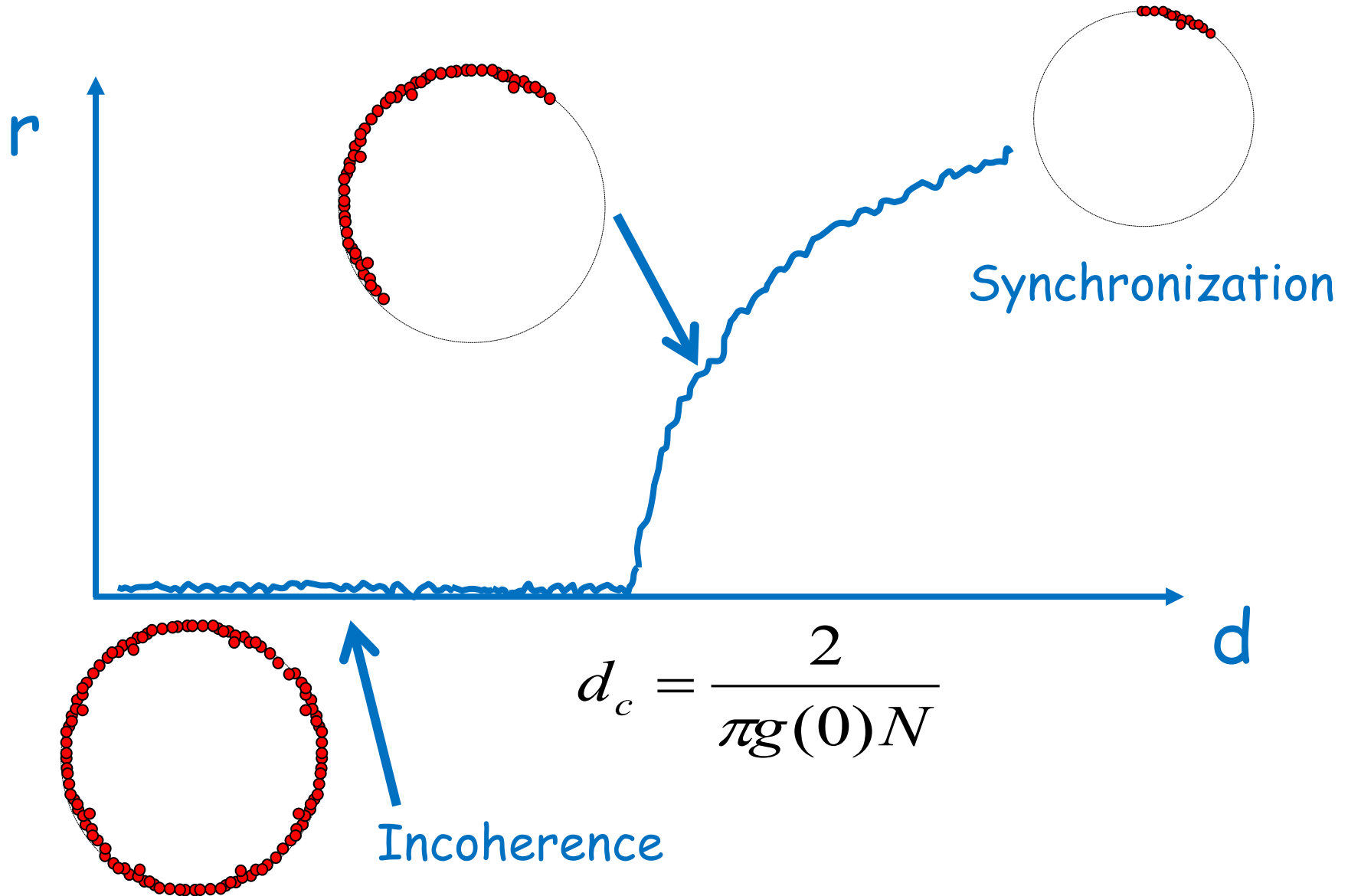


$r \approx 1$



$r \approx 0$

Path to synchrony in the Kuramoto model



Explosive

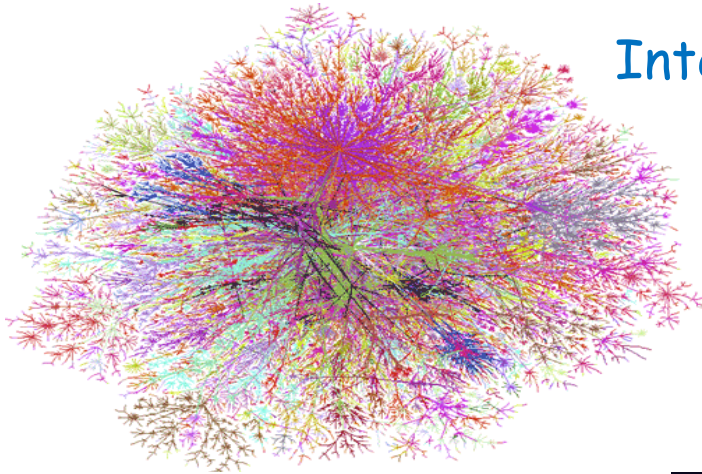
Synchronization

in

Complex Networks

Why complex networks?

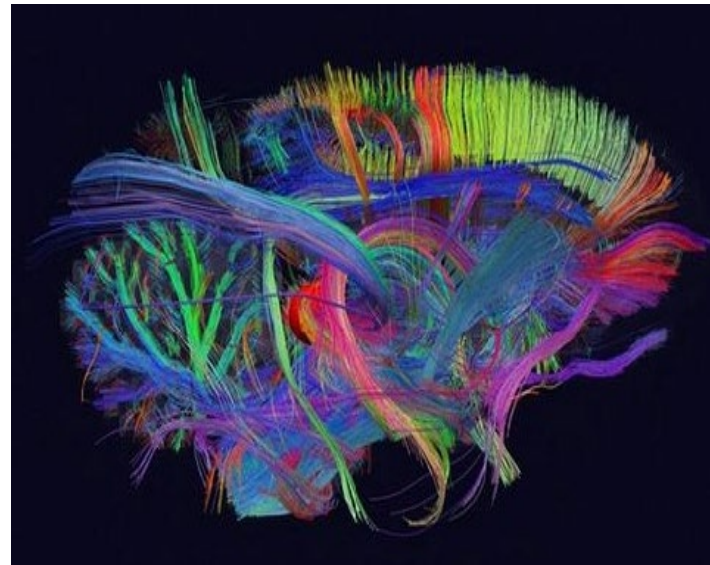
Most of the systems where synchronization is important are **complex networks**



Internet



Power grid



Brain



Next step: synchronization of Kuramoto units in complex networks



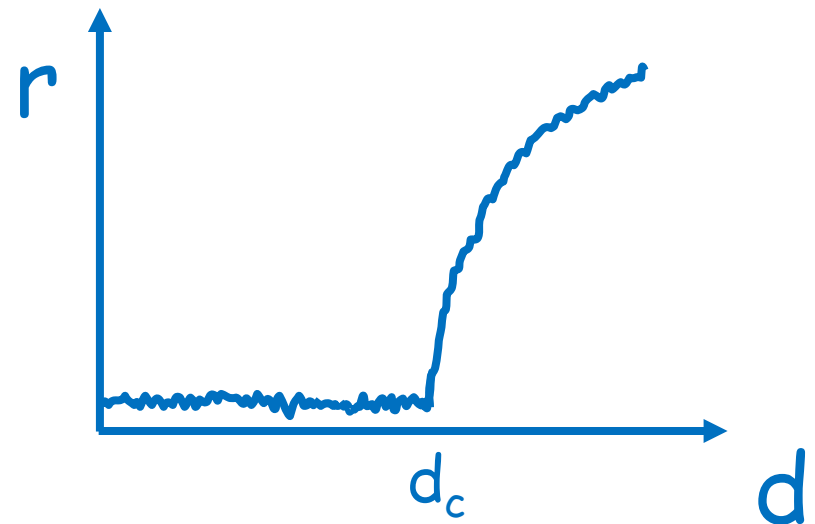
In order to study the effect of a network on the emergence of synchronization, we will maintain the simple phase dynamics, but will introduce a **complex network** in the problem

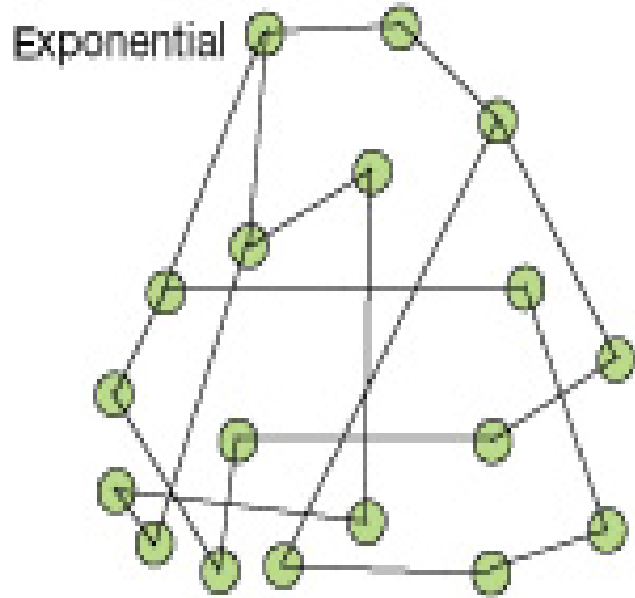
$$\dot{\theta}_n = \omega_n + d \sum_{m=1}^N A_{nm} \sin(\theta_m - \theta_n)$$

A red arrow points from the text 'complex network' to the adjacency matrix term A_{nm} in the equation.

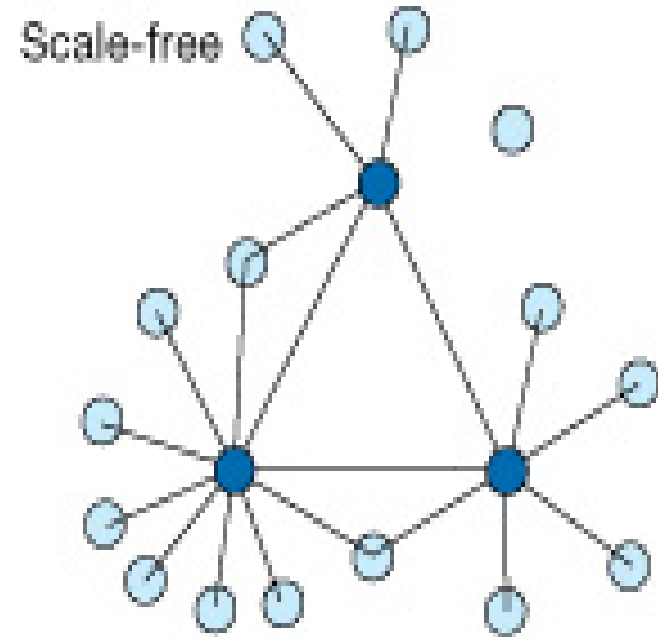
If m links to n $\leftrightarrow A_{nm} = 1$
(else $A_{nm} = 0$)

$$d_c = \frac{2}{\pi g(0)} \boxed{?}$$





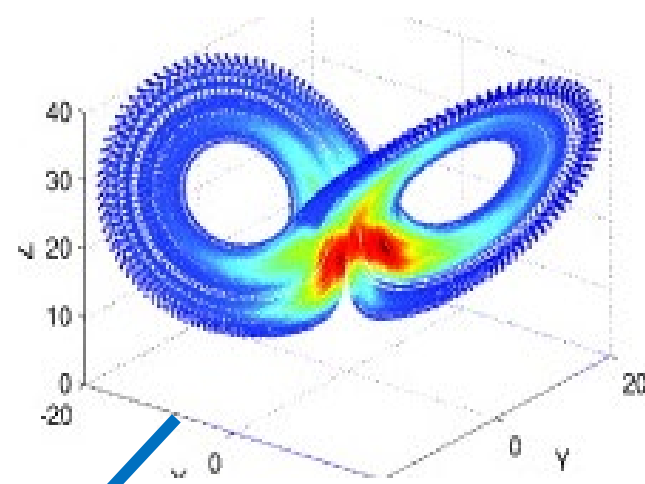
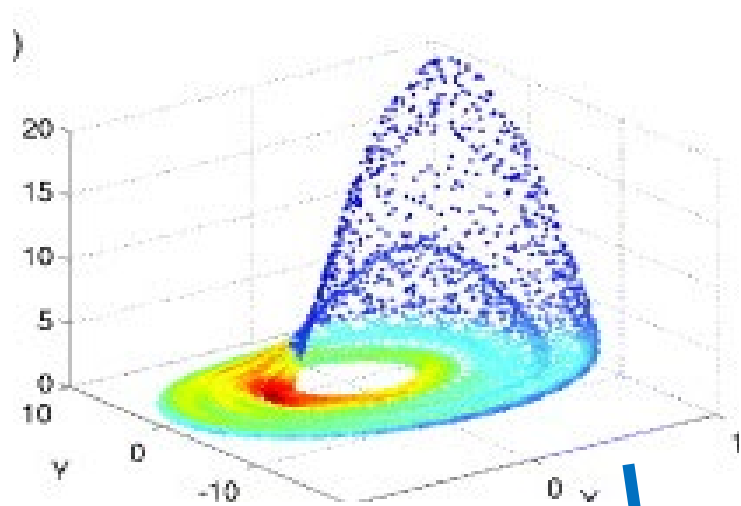
$$d_c > d_c$$



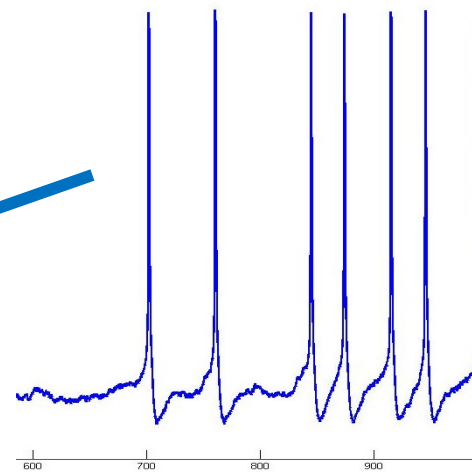
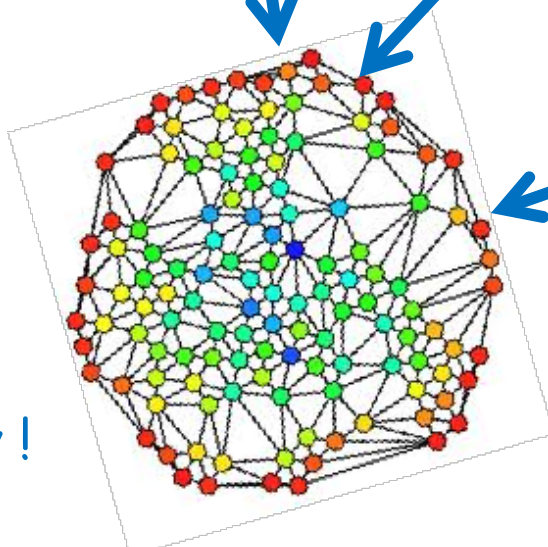
For a given number of nodes and connections,
heterogeneous networks synchronizes easier

Structure matters !

Complex dynamics in complex networks



Synchronization strongly depends on **both topology AND dynamics**



...in a way too long for telling you today !

Explosive

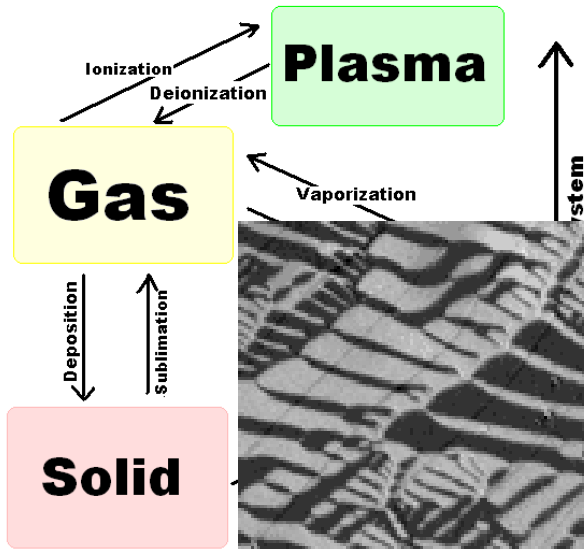
Synchronization

in

Complex Networks

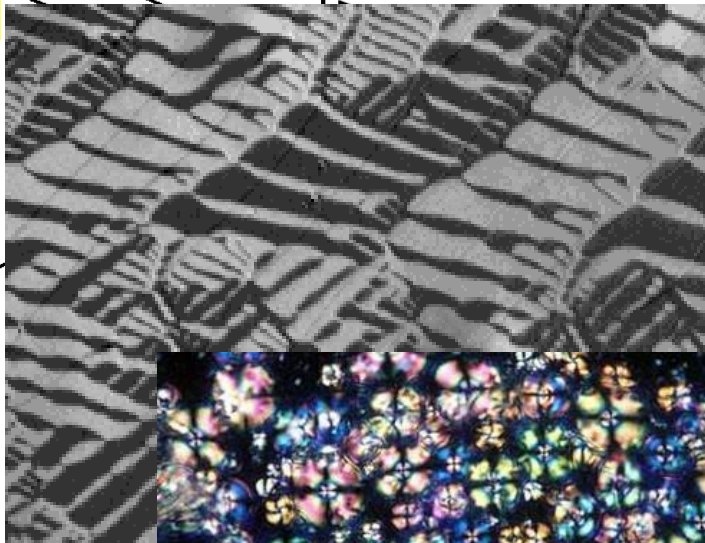
Phase transitions

Changes between different states of organization in a system

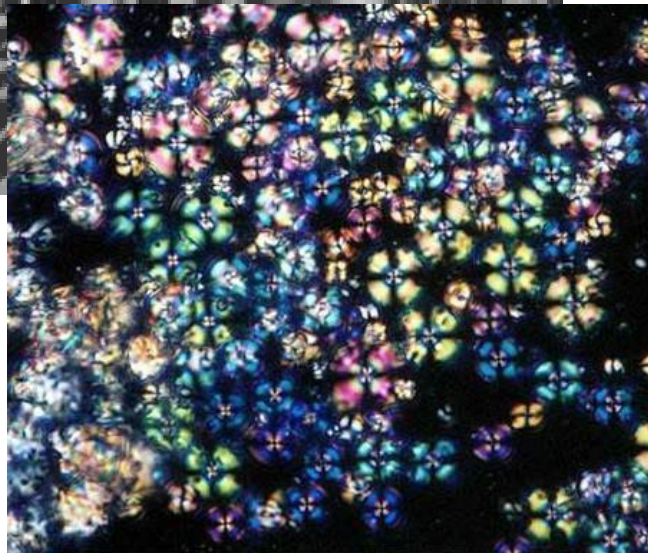


States of matter

Magnetization

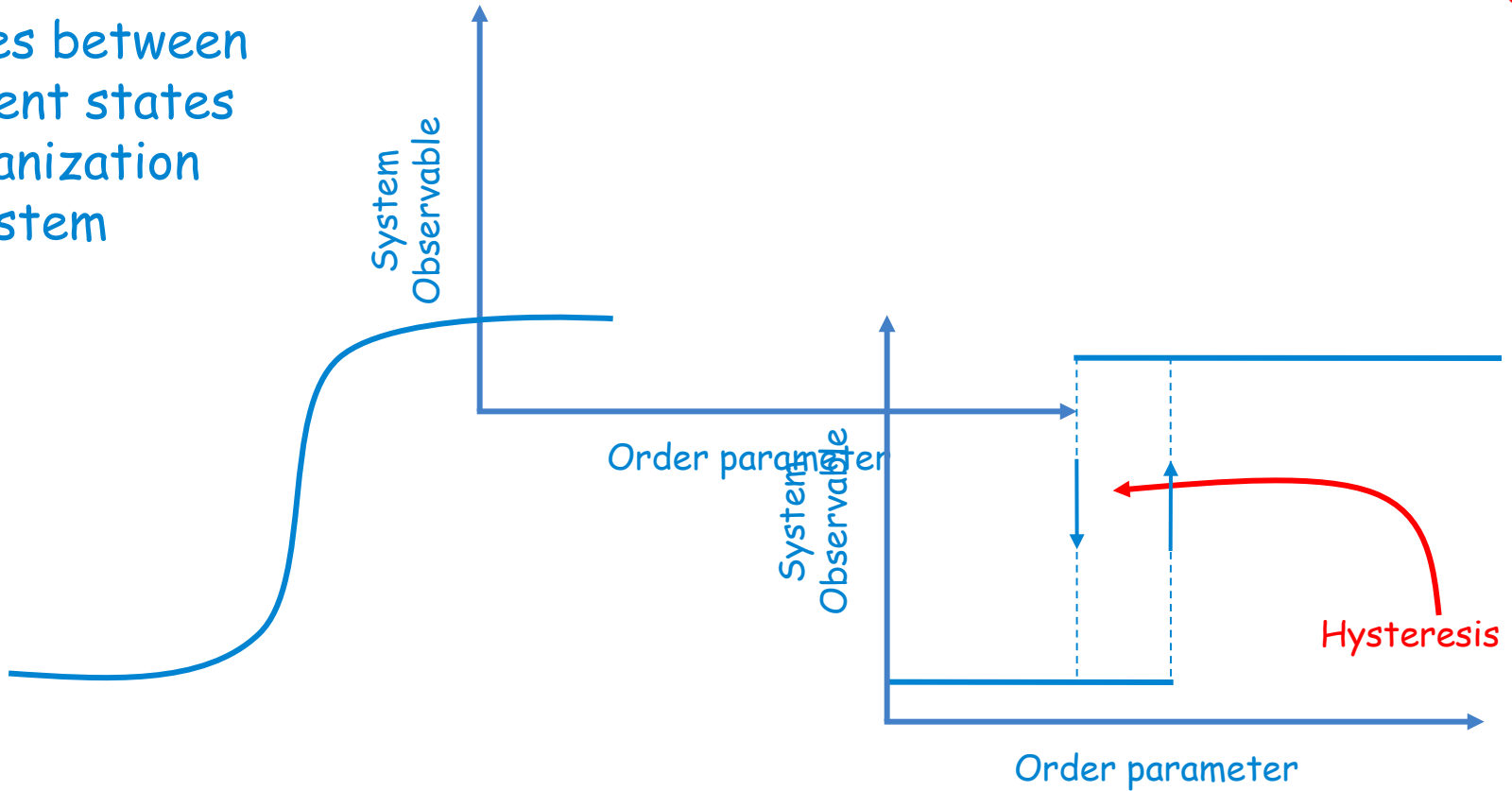


Liquid cristal nematic state



Remembering phase transitions

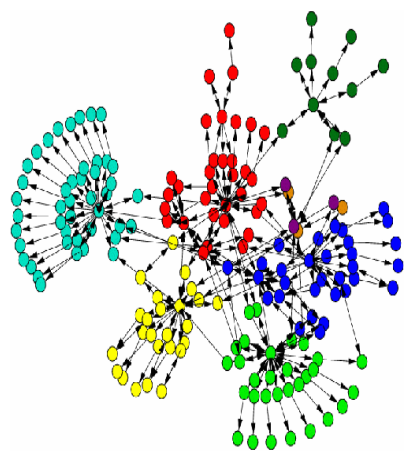
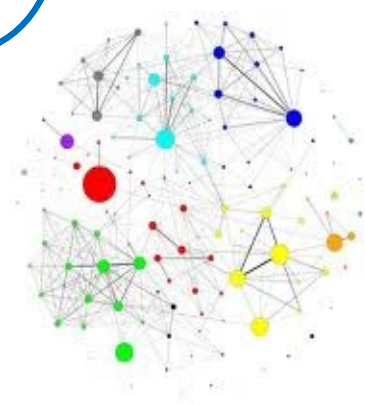
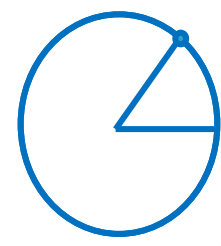
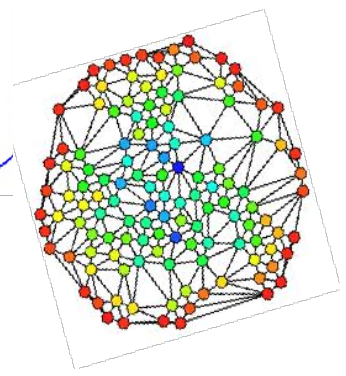
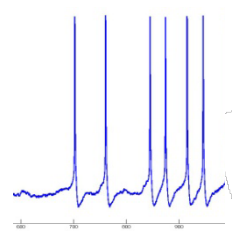
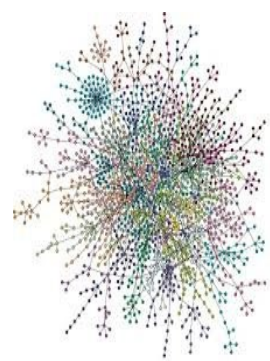
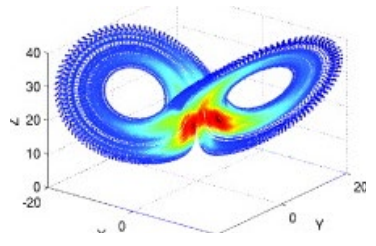
Changes between different states of organization in a system



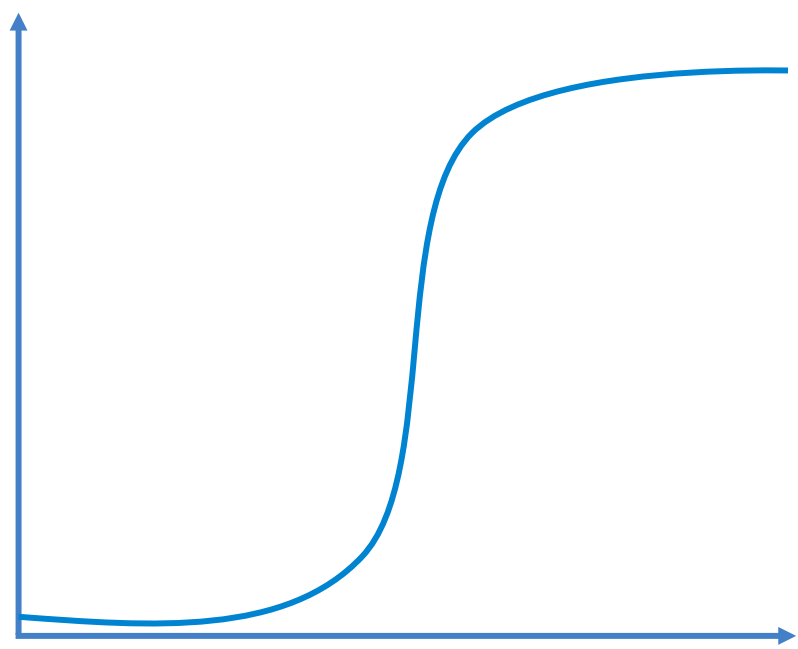
Second order PT

First order PT

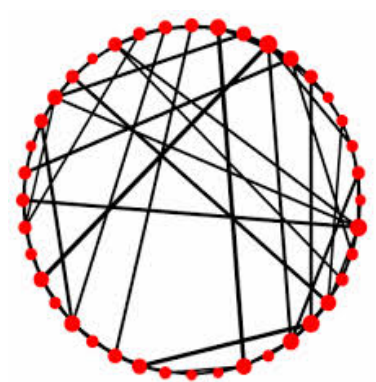
Transitions to synchrony in complex networks



Synchronization



Coupling strength



Case 1: Kuramoto model with a special frequency distribution



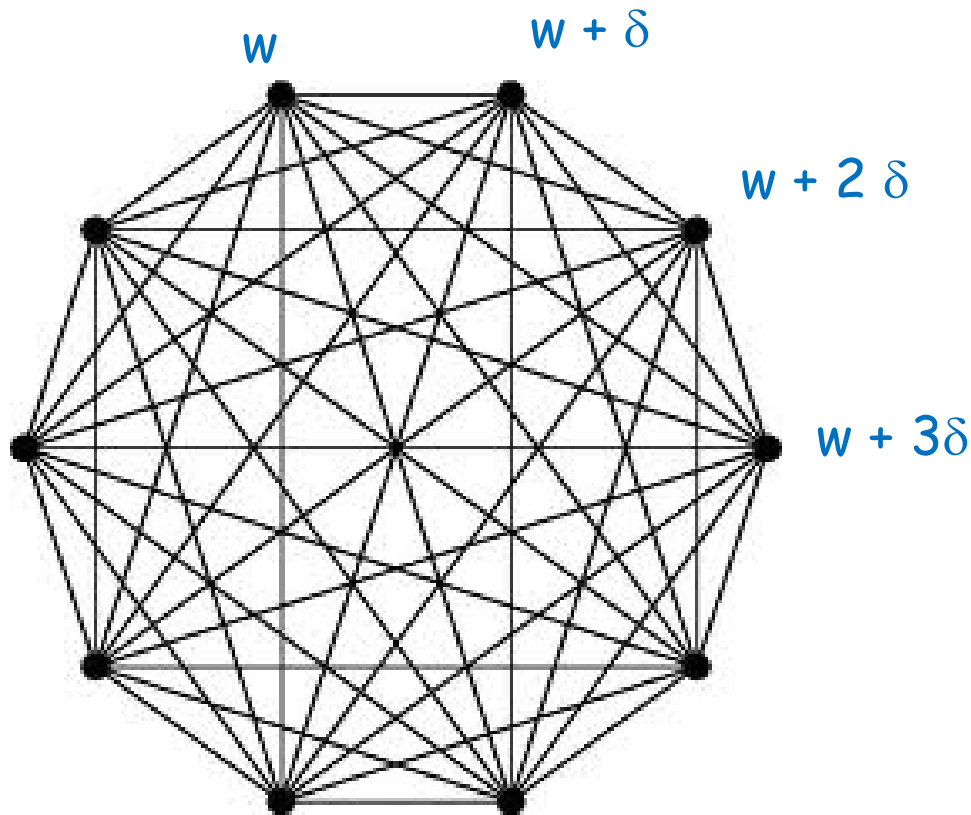
PHYSICAL REVIEW E 72, 046211 (2005)

Thermodynamic limit of the first-order phase transition in the Kuramoto model

Diego Pazó*

Max-Planck-Institut für Physik komplexer Systeme, Nöthnitzer Straße 38, 01187 Dresden, Germany

Full Kuramoto model with *equispaced* frequencies



$$\omega_j = -\gamma + \frac{\gamma}{N} (2j - 1)$$

Case 2: SF + degree-frequency correlation

PRL 106, 128701 (2011)

PHYSICAL REVIEW LETTERS

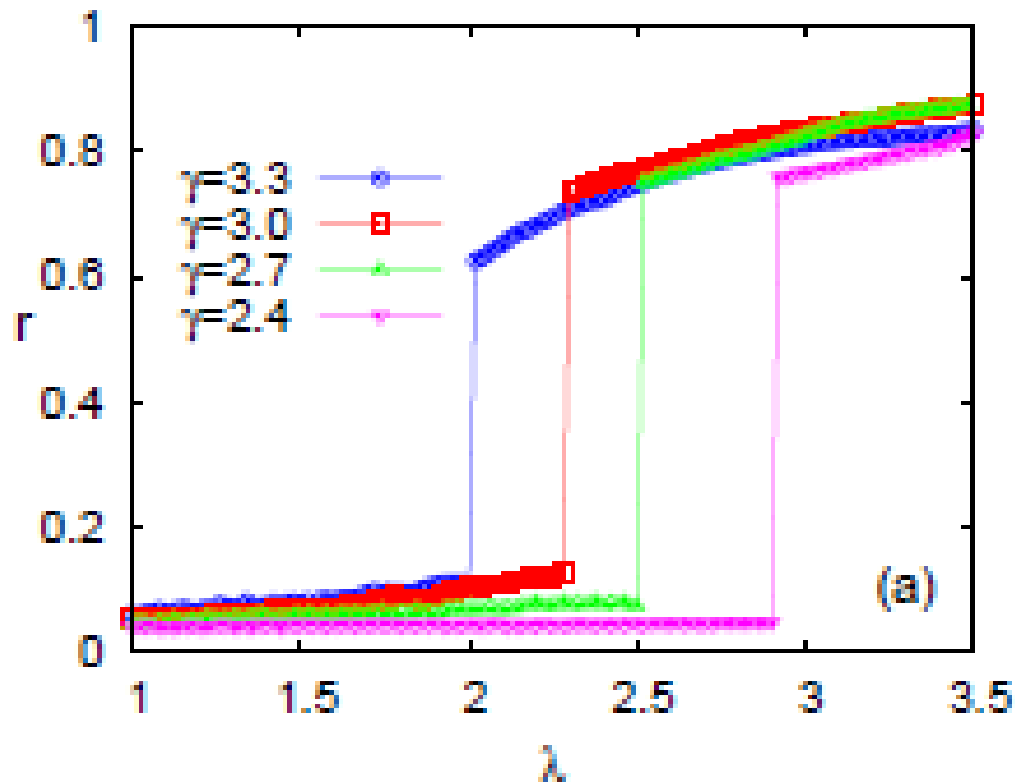
week ending
25 MARCH 2011



Explosive Synchronization Transitions in Scale-Free Networks

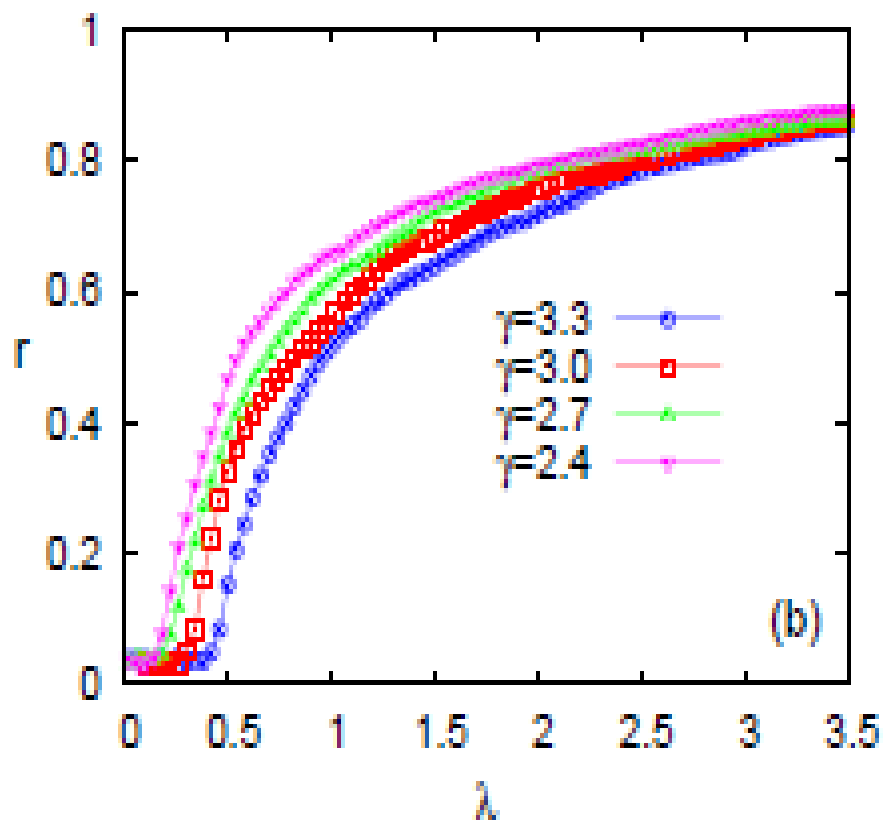
Jesús Gómez-Gardeñes,^{1,2,*} Sergio Gómez,³ Alex Arenas,^{2,3} and Yamir Moreno^{2,4}

SF networks of Kuramoto oscillators where ω_i proportional to k_i

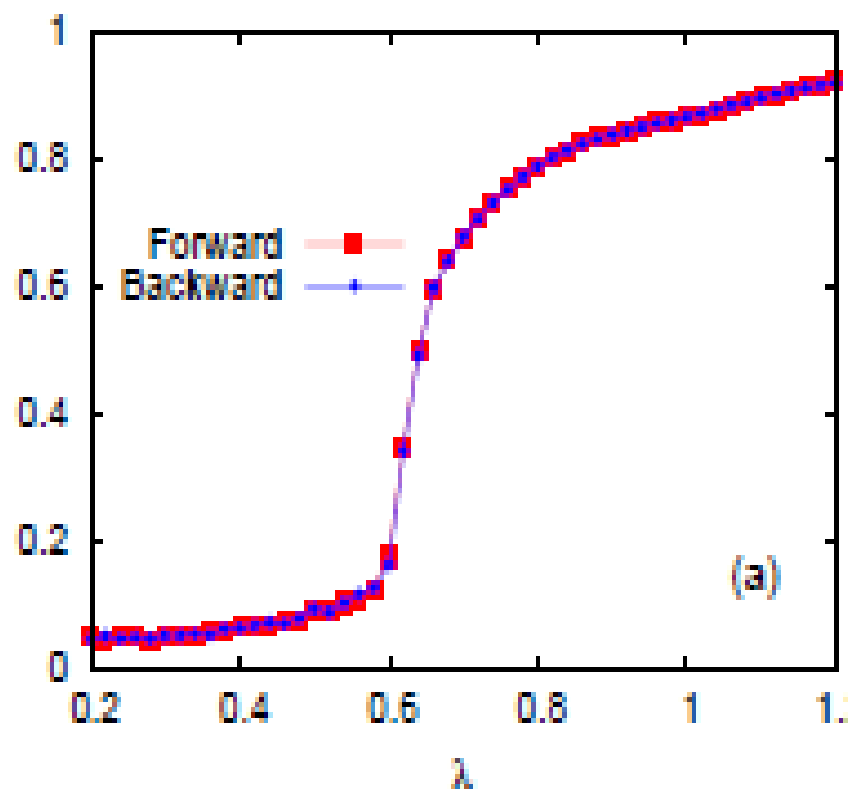


Case 2: SF network + degree-frequency correlation

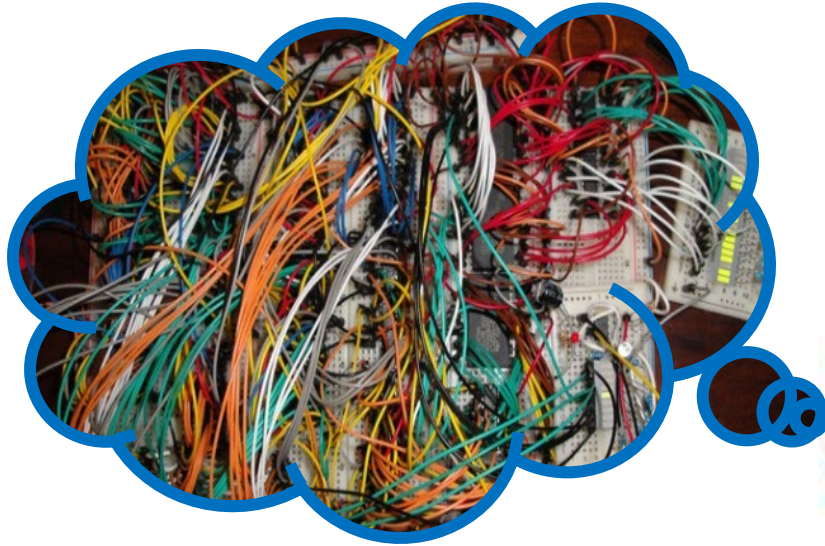
SF but
uncorrelated ω_i, k_i



Correlated ω_i, k_i
but random network



(What the coffe-breaks are useful for...)



Wouldn't be *cool* if this could be done
EXPERIMENTALLY ?

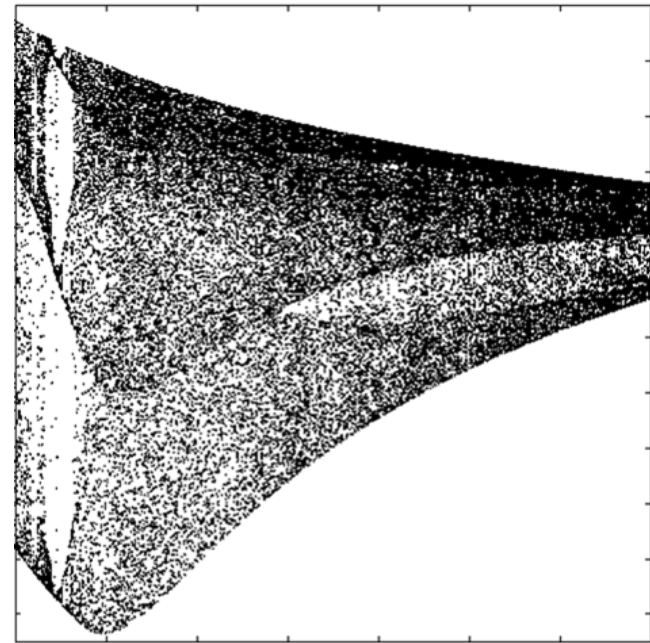
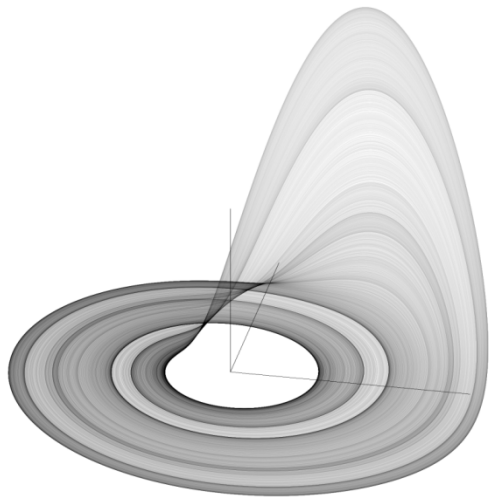
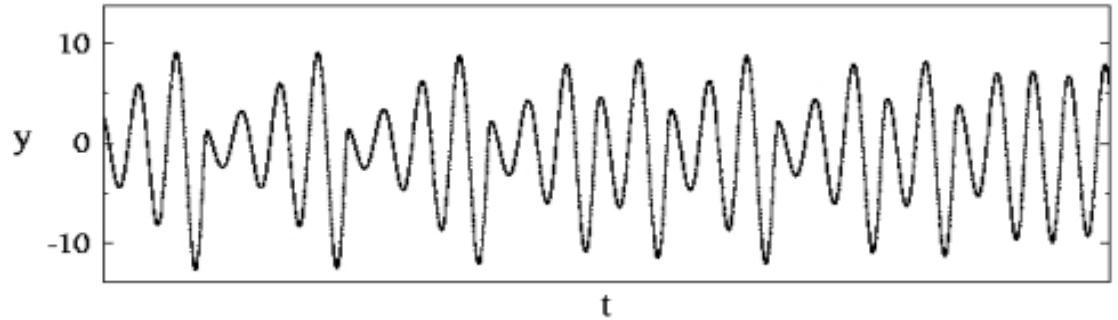
The model: Rössler oscillator

Even more fun: Can it be done experimentally

$$\dot{x} = -\omega y - z$$

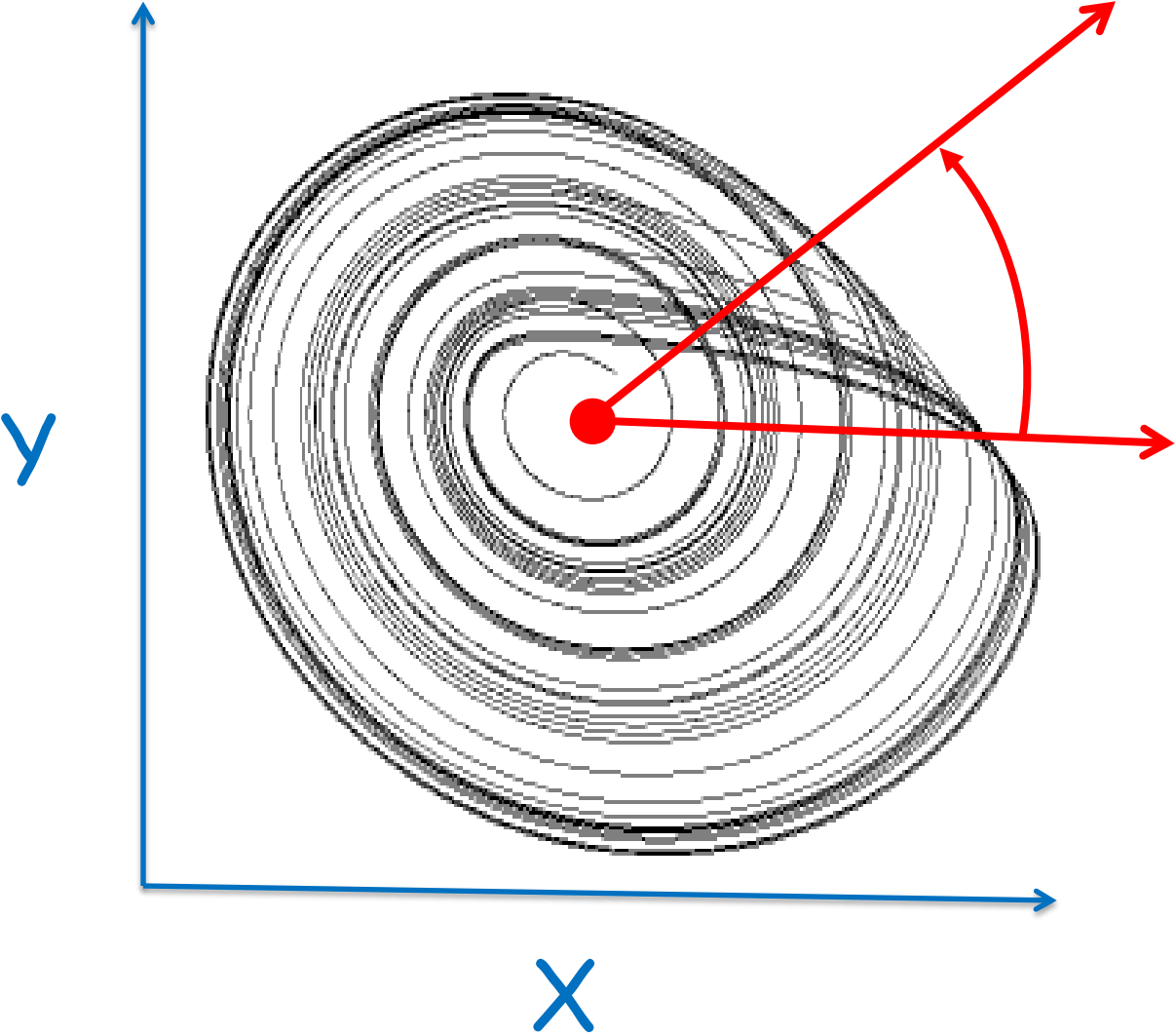
$$\dot{y} = \omega x + \alpha y$$

$$\dot{z} = b + z(x - c)$$



a

Phase and amplitude in the chaotic Rössler oscillator



Phase

$$\theta = \text{artag}\left(\frac{y}{x}\right)$$

The *actual* model: piecewise Rössler oscillator

$$\dot{x}_i = -\alpha_i (\Gamma x_i + \beta y_i + \lambda z_i) + d \sum_{j=1}^N a_{ij} (x_j - x_i)$$

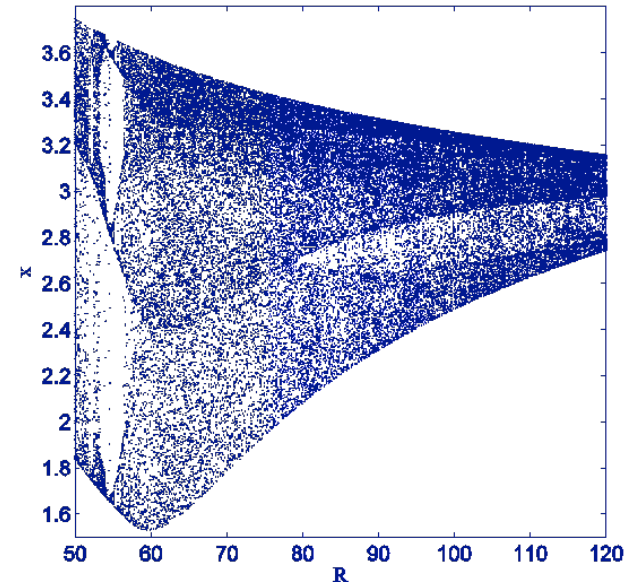
$$\dot{y}_i = -\alpha_i \left(-x_i + \left(m - \frac{n}{R} \right) y_i \right)$$

$$\dot{z}_i = -\alpha_i \left(-\boxed{g(x_i)} + \lambda z_i \right)$$

piecewise
part

Frequency control in a very large range

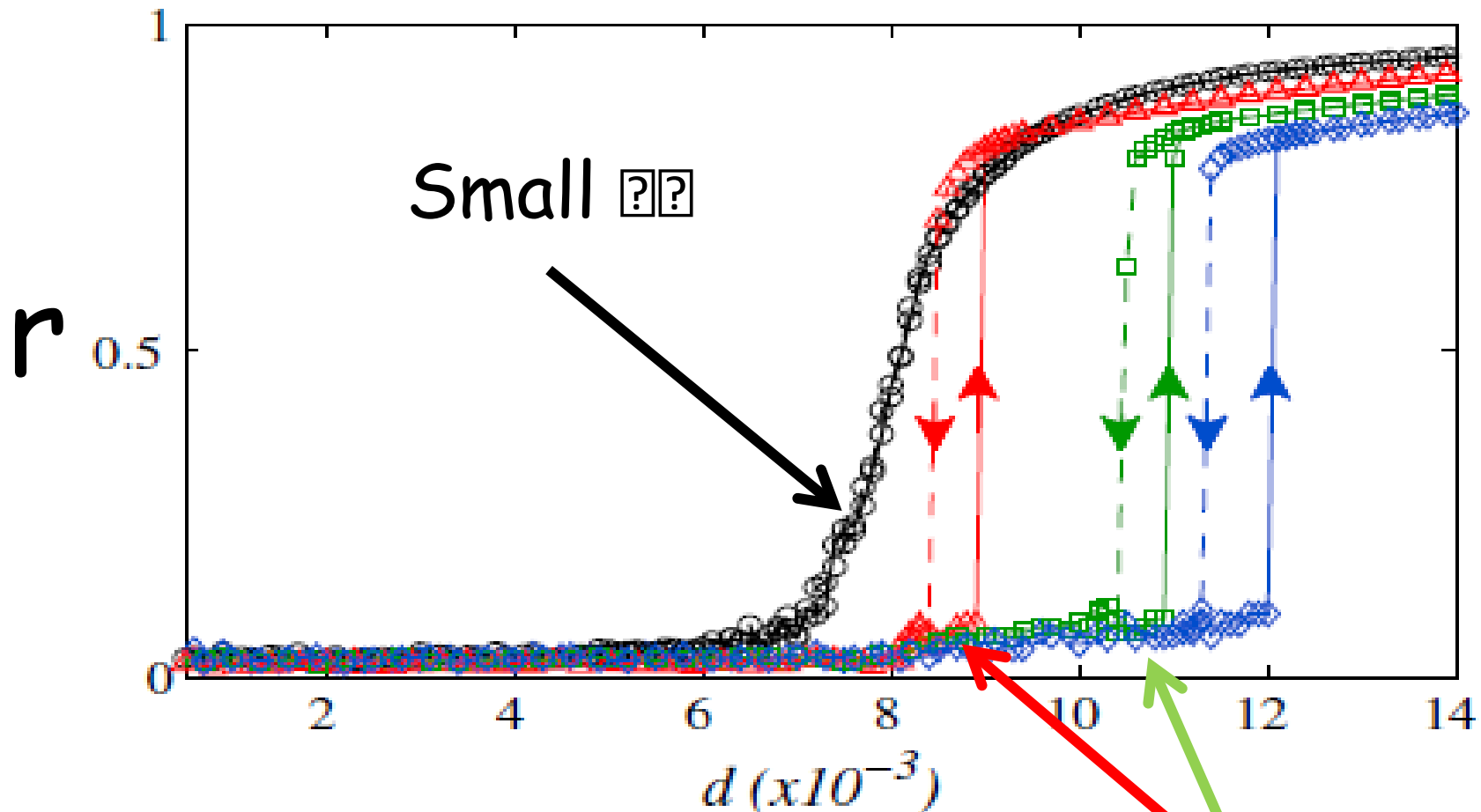
$$\alpha_i = \alpha \left(1 + \Delta\alpha \frac{k_i - 1}{N} \right)$$



Dynamical state
control

Simulation: explosive phase synchronization

SF network, $N=1000$, several values of $\Delta\alpha$



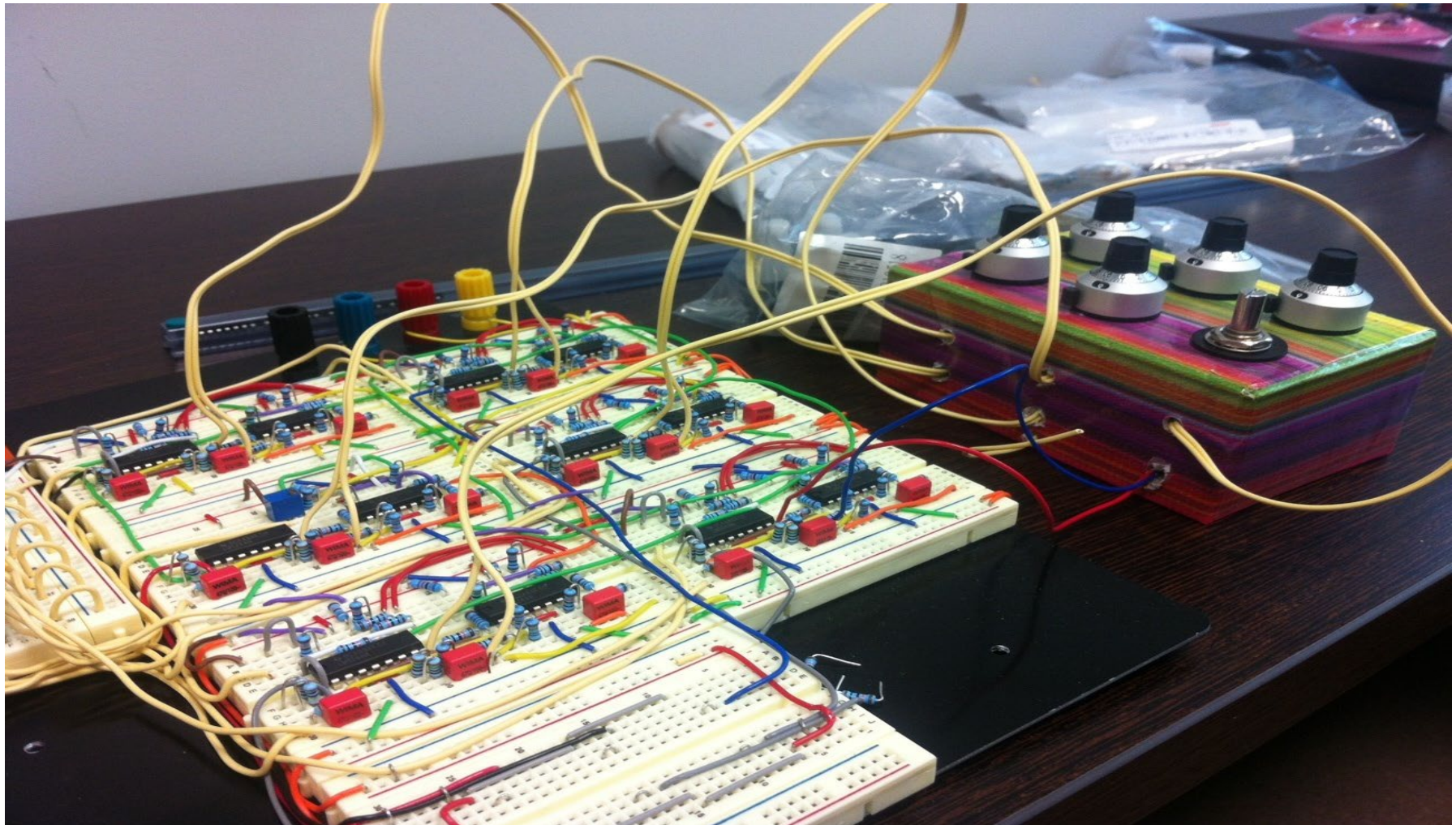
$$\alpha_i = \alpha \left(1 + \Delta\alpha \frac{k_i - 1}{N} \right)$$

Large $\Delta\alpha$

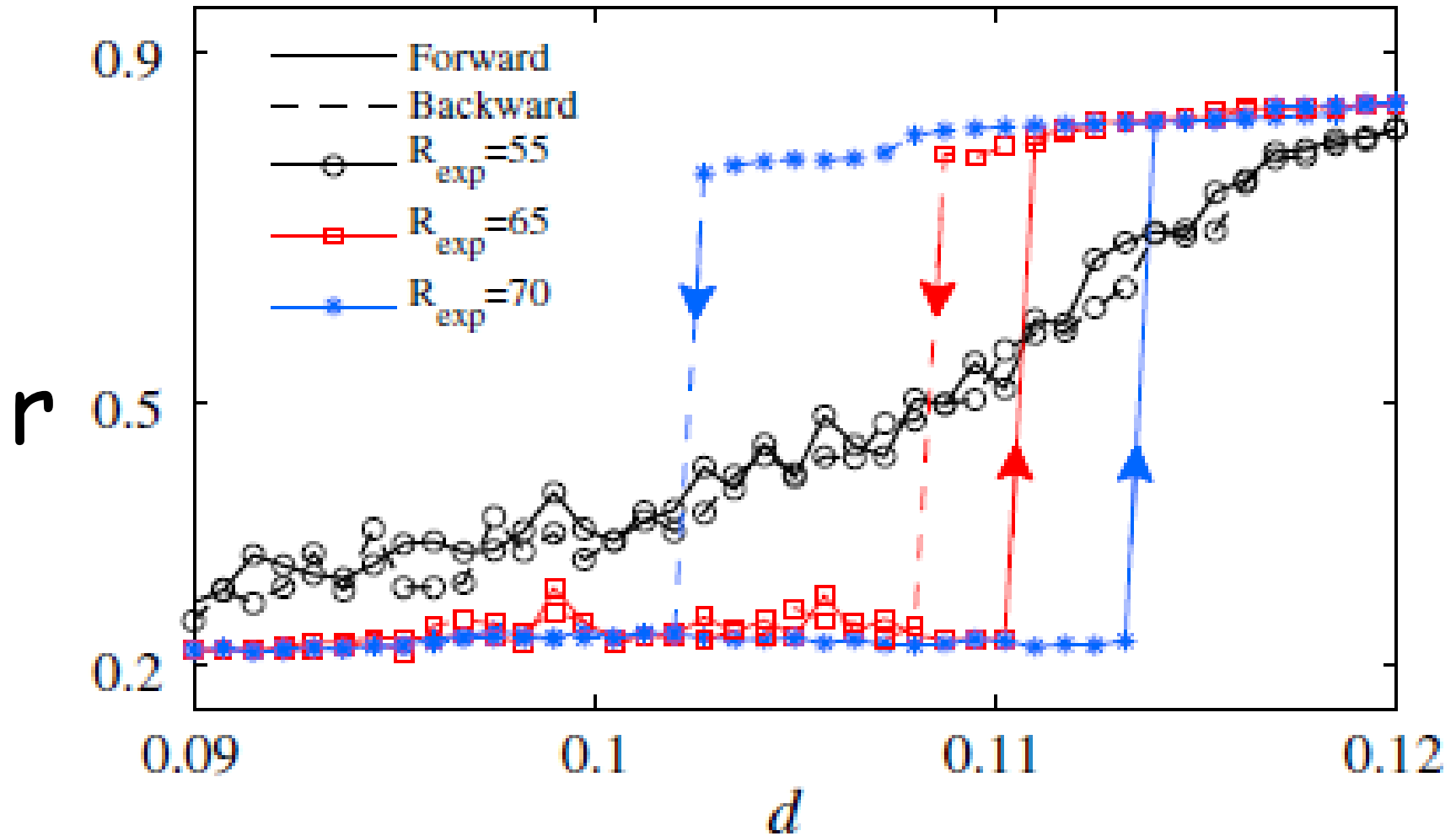
The experiment: star network

$N = 5 + 1$ with common R parameter \rightarrow same dynamical state

- Fast node N_1 $\tau_1 = 3333$ Hz
- Slow nodes $N_2 \dots N_6$ $\tau_i = 2240 \tau_{200}$ Hz

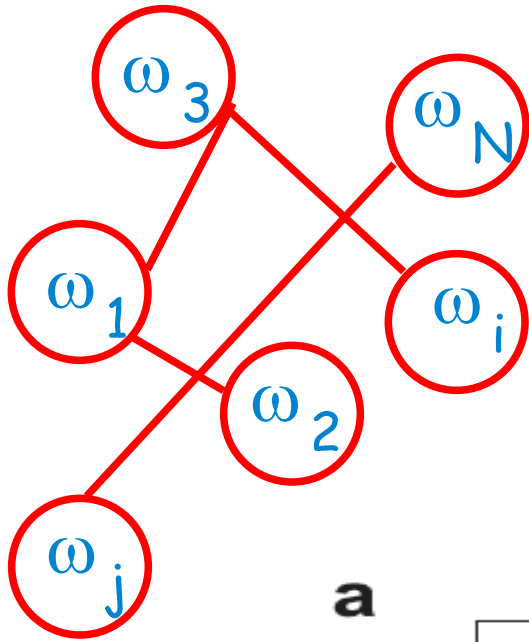


Experimental explosive synchronization



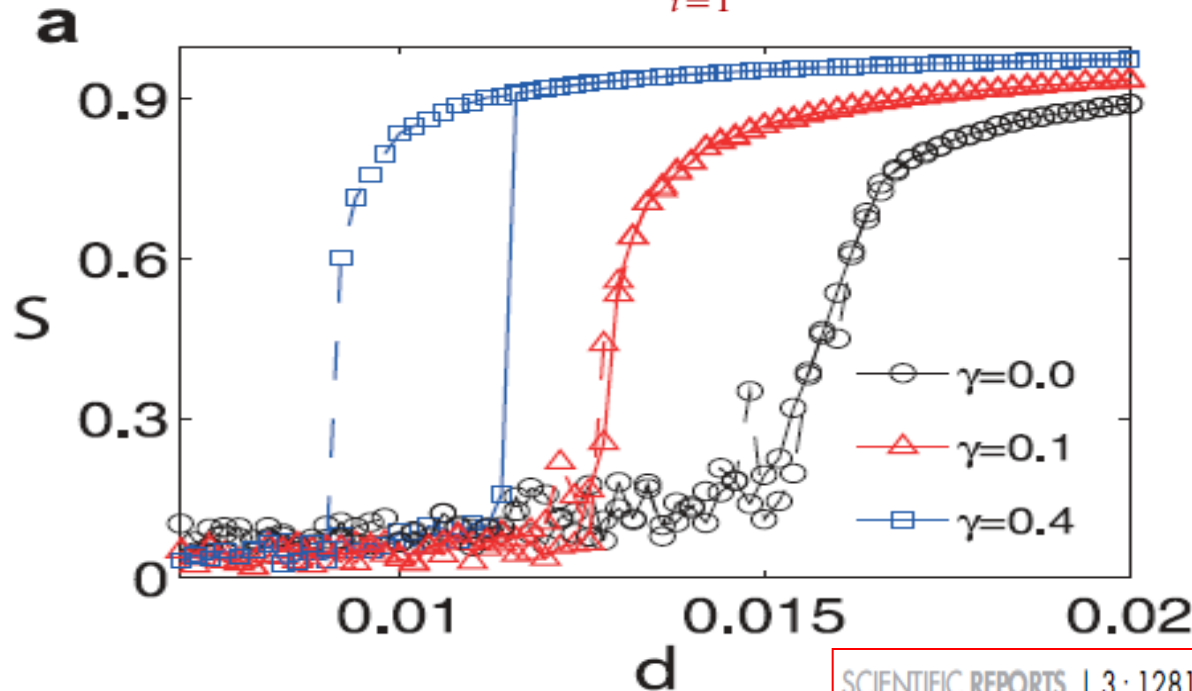
First experimental evidence of first order synchronization

Making you network to explode I : Gap method



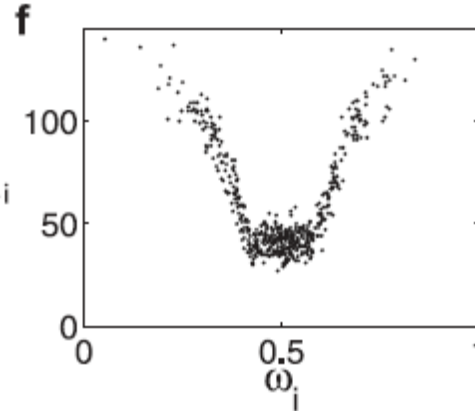
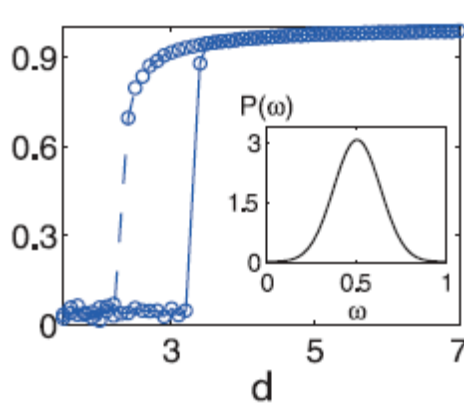
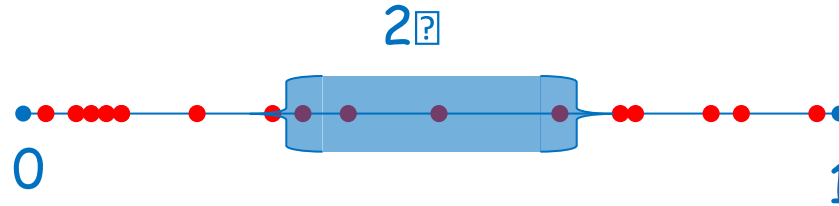
- Distribute frequencies (valid for any $g(w)$)
- Pick a random pair i, j
- Only if $|\omega_i - \omega_j| > \gamma \rightarrow a_{ij} = 1$
- Continue up to construct target network ($\langle k \rangle ?$)

$$\frac{d\phi_i}{dt} = \omega_i + d \sum_{i=1}^N a_{ij} \sin(\phi_j - \phi_i),$$

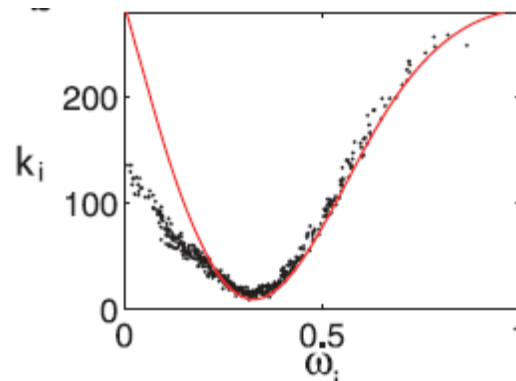
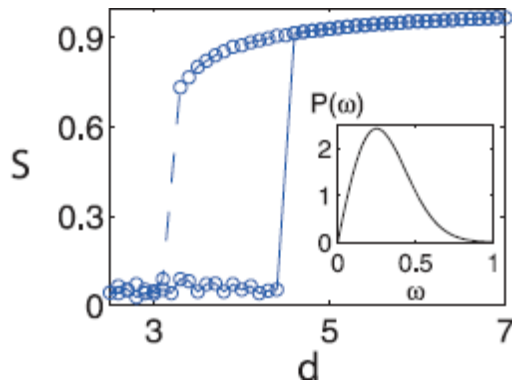


Gap method: emergent k - ω correlation

Probability of node i of being linked depends on $\omega_i \rightarrow g(\omega)$



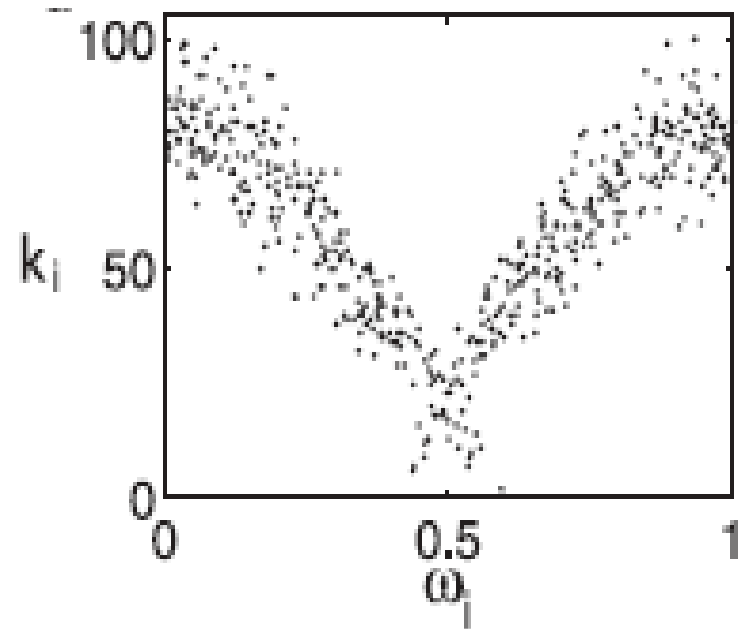
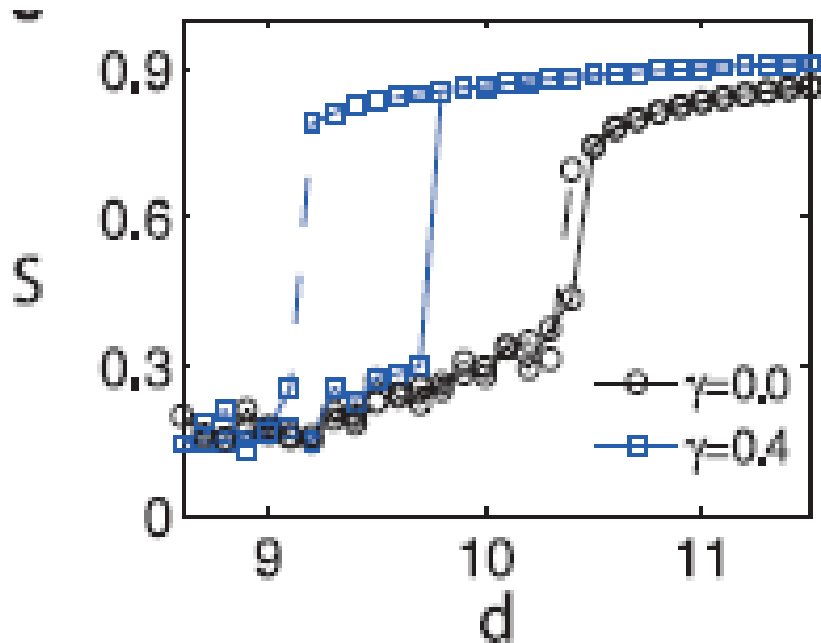
... but work equally for regular random networks (all nodes with same k)



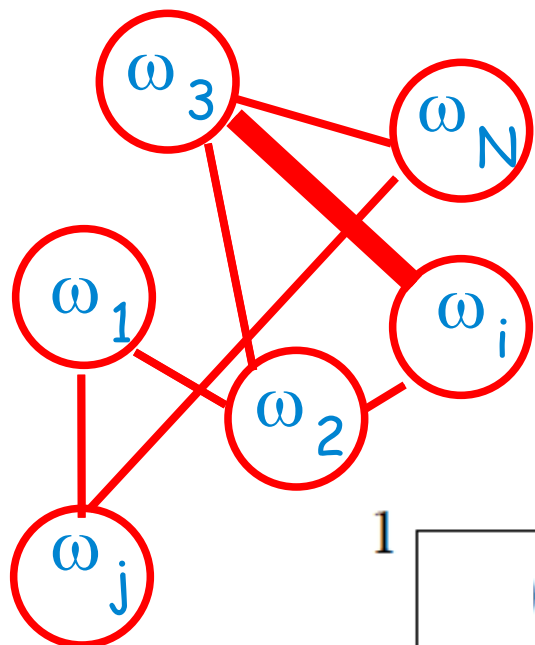
k - ω correlation can emerge, *but is not a condition*

Also works for weaker rules as **neighbourhood averaged gap**:

$$\begin{aligned} & \left| \omega_i - \langle \omega_j \rangle \right| > \gamma \\ & \left| \omega_j - \langle \omega_i \rangle \right| > \gamma \end{aligned} \rightarrow a_{ij}=1$$

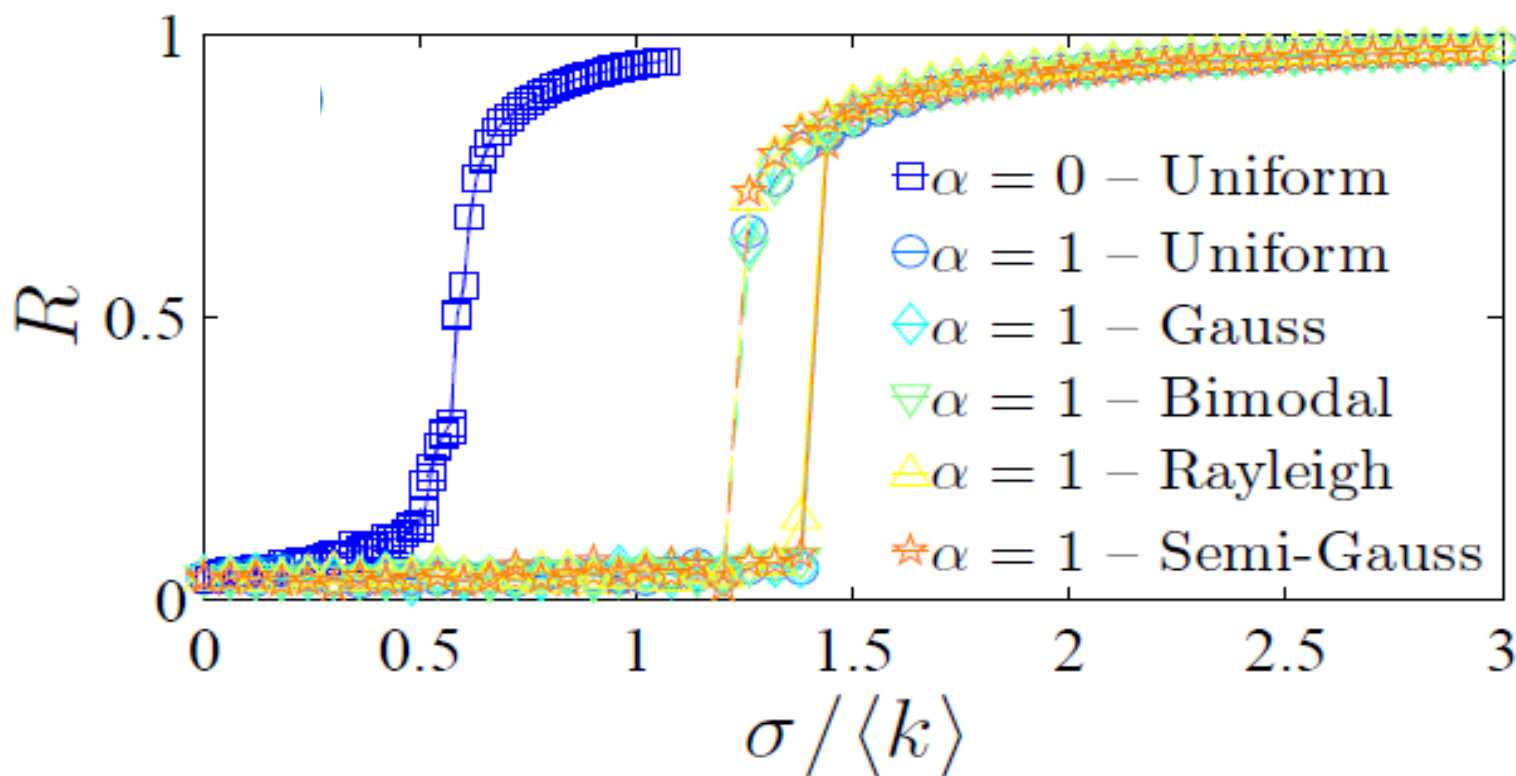


Making your network to explode II : weighting method

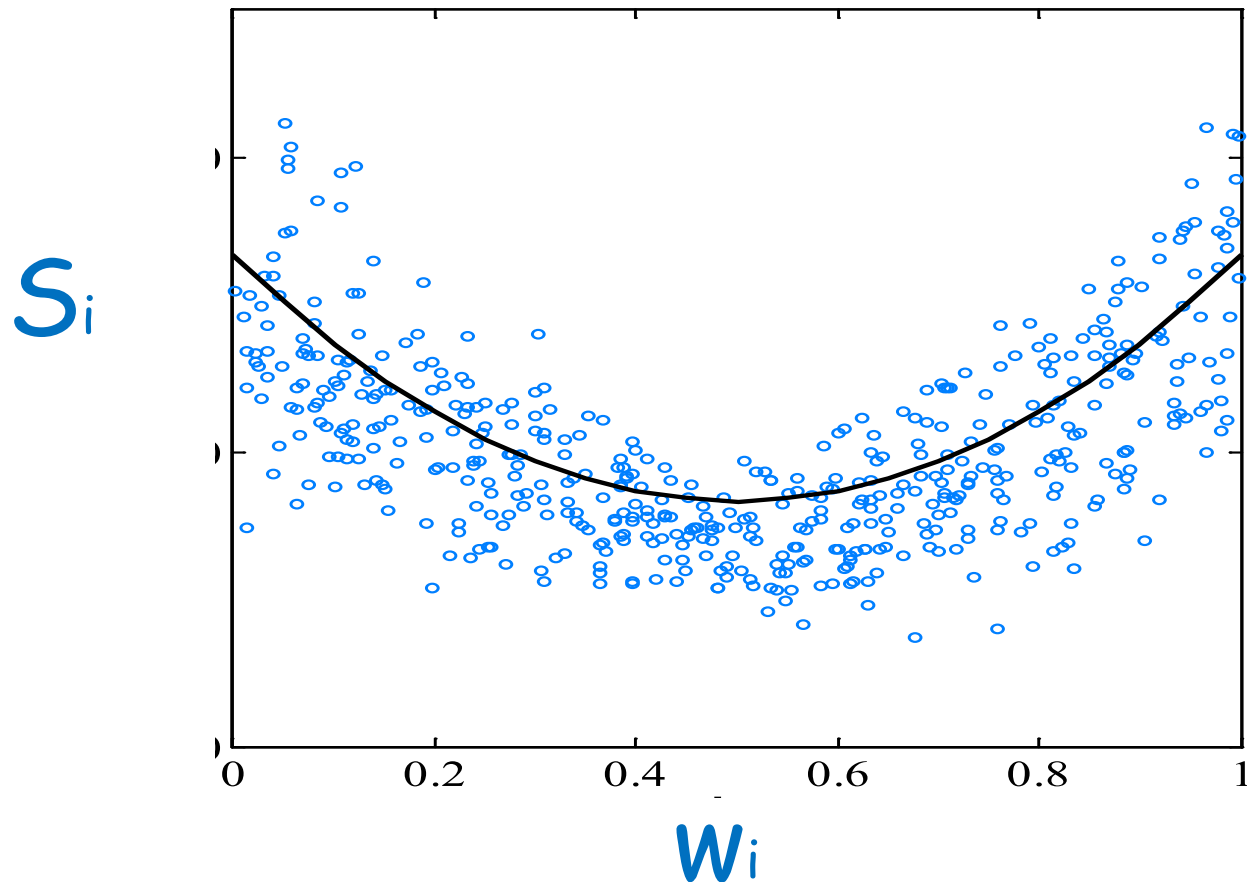


$$\frac{d\theta_i}{dt} = \omega_i + \frac{\sigma}{\langle k \rangle} \sum_{i=1}^N a_{ij} \sin(\theta_j - \theta_i),$$

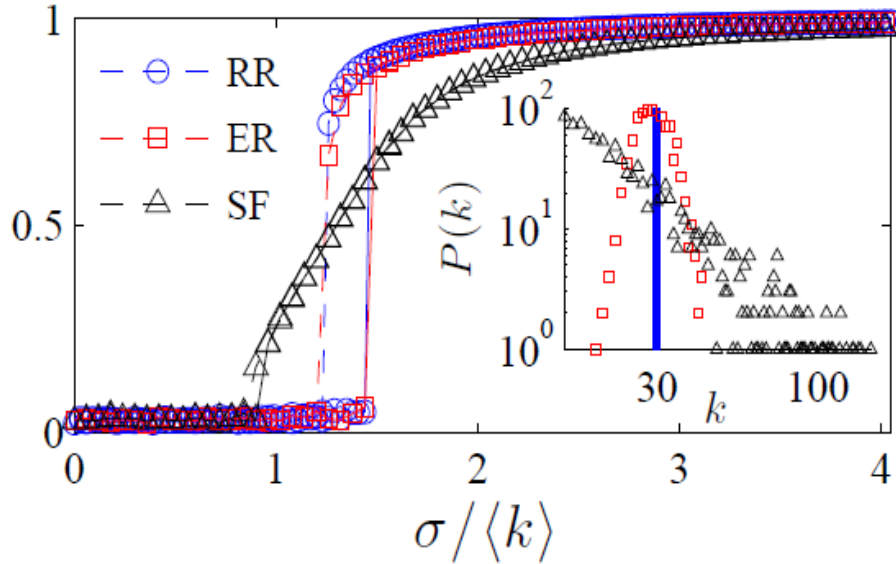
$$\Omega_{ij}^\alpha = a_{ij} |\omega_i - \omega_j|^\alpha$$



Node strength $S_i = \sum_j \Omega_{ij}$



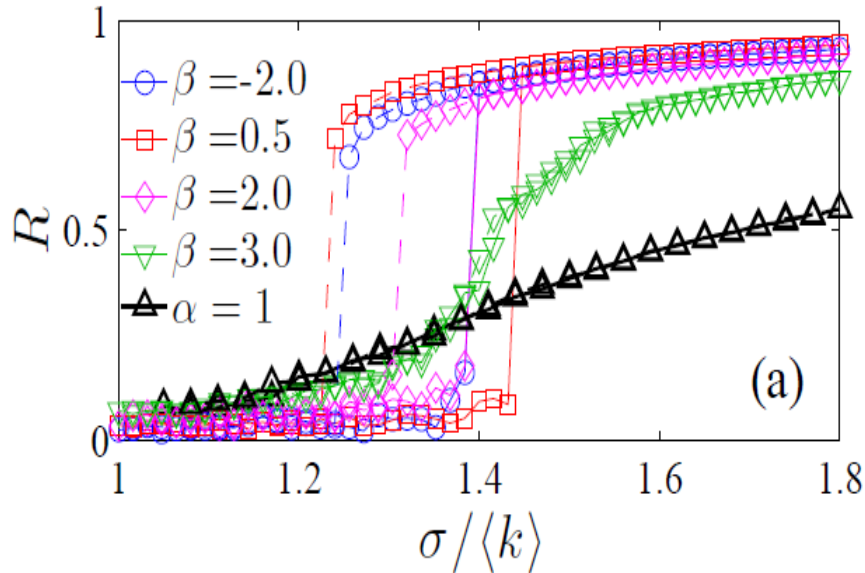
ES and the heterogeneity paradox



Heterogeneous networks need a *detuning/topology* weighting:

$$\tilde{\Omega}_{ij} = a_{ij} |\omega_i - \omega_j| \frac{l_{ij}^\beta}{\sum_{j \in \mathcal{N}_i} l_{ij}^\beta}$$

l_{ij} edge betweenness of a_{ij}



Explosive synchronization for $\beta > 0$
(maximum hysteresis width $\beta = 0.5$)

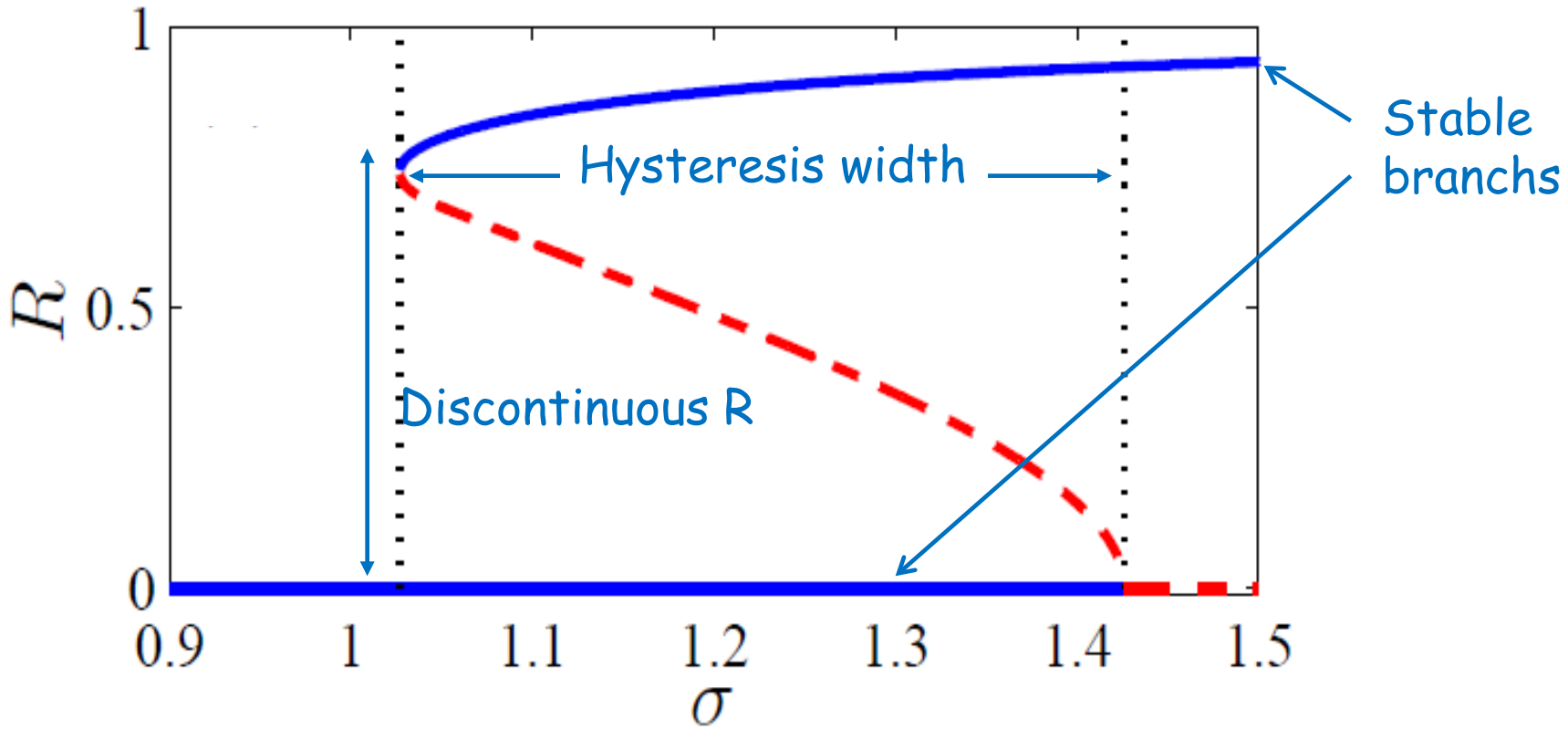
In the thermodynamic limit

$$\dot{\theta}_i = \omega_i + \frac{\sigma}{N} \sum_{j=1}^N \Omega_{ij} \sin(\theta_j - \theta_i),$$

Co-rotating frame phases

$$\omega = \sigma A_\omega \sin(\theta_\omega - \phi_\omega).$$

where $A_\omega \sin \phi_\omega = \int g(x) |w - x| \sin \theta(x) dx$ depends on \square



Are correlations necessary? Answer is NO!



Explosive Synchronization in adaptive and multi-layer networks

X. Zhang, S. Boccaletti, S. Guan, Z. Liu, Phys. Rev. Lett. **114**, 038701 (2015)

$$\Theta'_i = \omega_i + \lambda \alpha_i \sum_{j=1}^N A_{ij} \sin(\theta_j - \theta_i)$$

See.....pdf.....

What after? The Bellerophon states



Coexistence of quantized, time dependent clusters in globally coupled oscillators

H. Bi, X. Hu, S. Boccaletti, X. Wang, Y. Zou, Z. Liu and S Guan,
Phys. Rev. Lett. **117**, 204101 (2016)

Synchronization

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Complex networks

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Explosive synchronization in Complex Networks

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- * X. Zhang, S. Boccaletti, S. Guan, Z. Liu, Phys. Rev. Lett. **114**, 038701 (2015)
- * S. Boccaletti et al., "Explosive transitions in complex networks' structure and dynamics: Percolation and synchronization", Phys. Rep. **660**, 1 (2016)

Thank you

(for your patience)

and.....

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CHAOS, SOLITONS & FRACTALS

EiC: Maurice Courbage and S.B.